

Equation of state of strong interactions: Expansion schemes and Critical point

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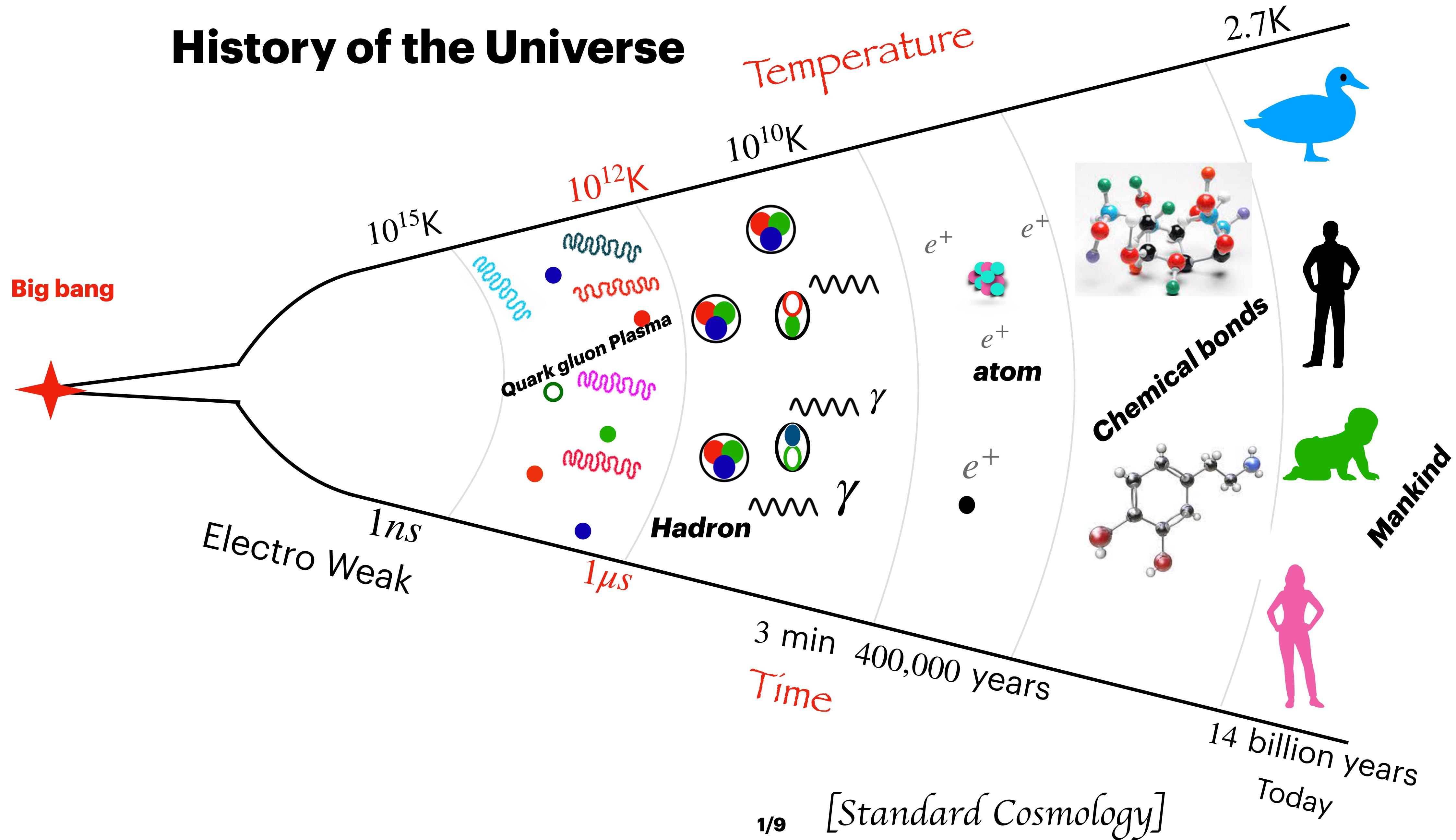
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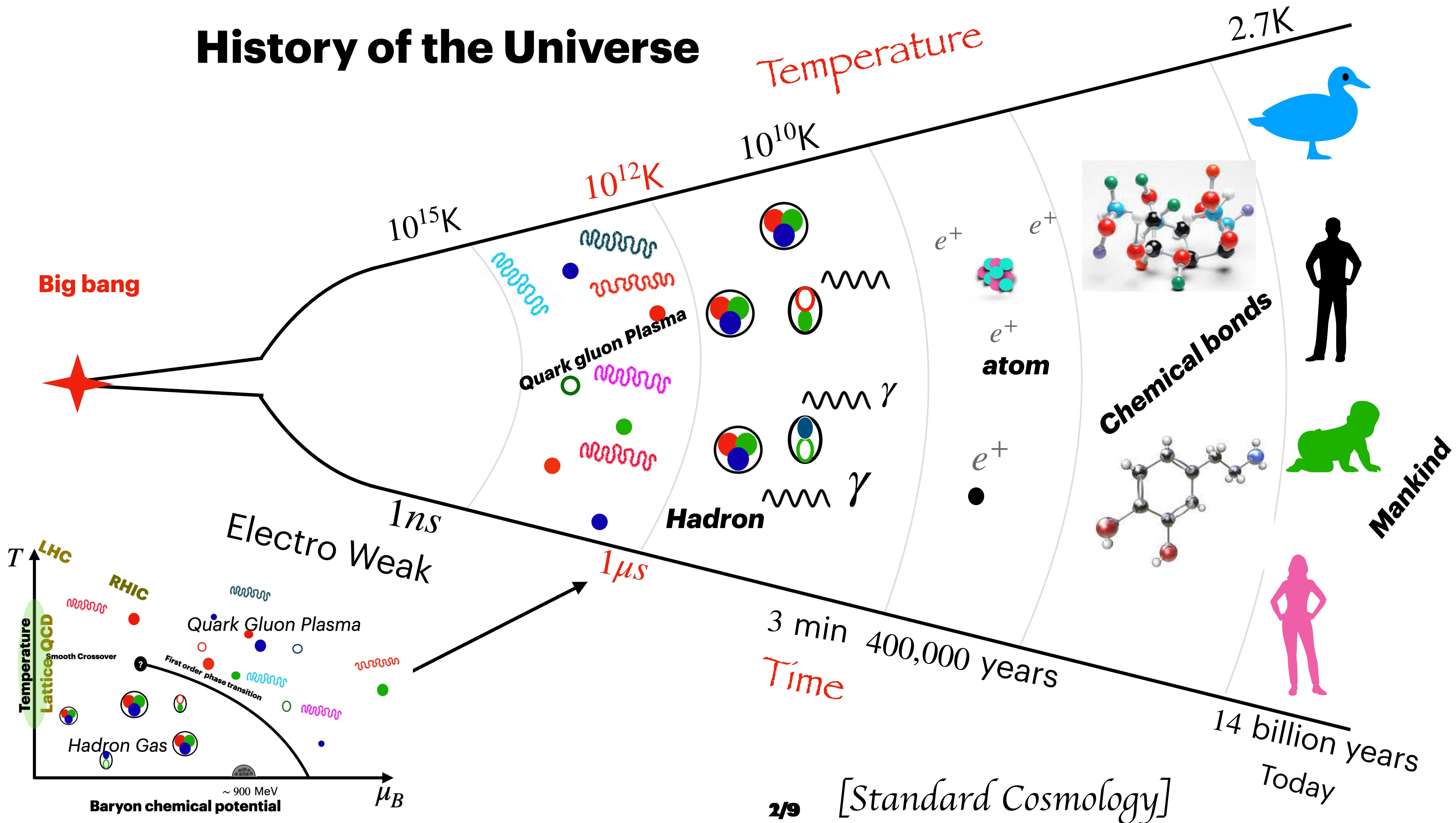
2024 NSBP - NSHP Joint Conference

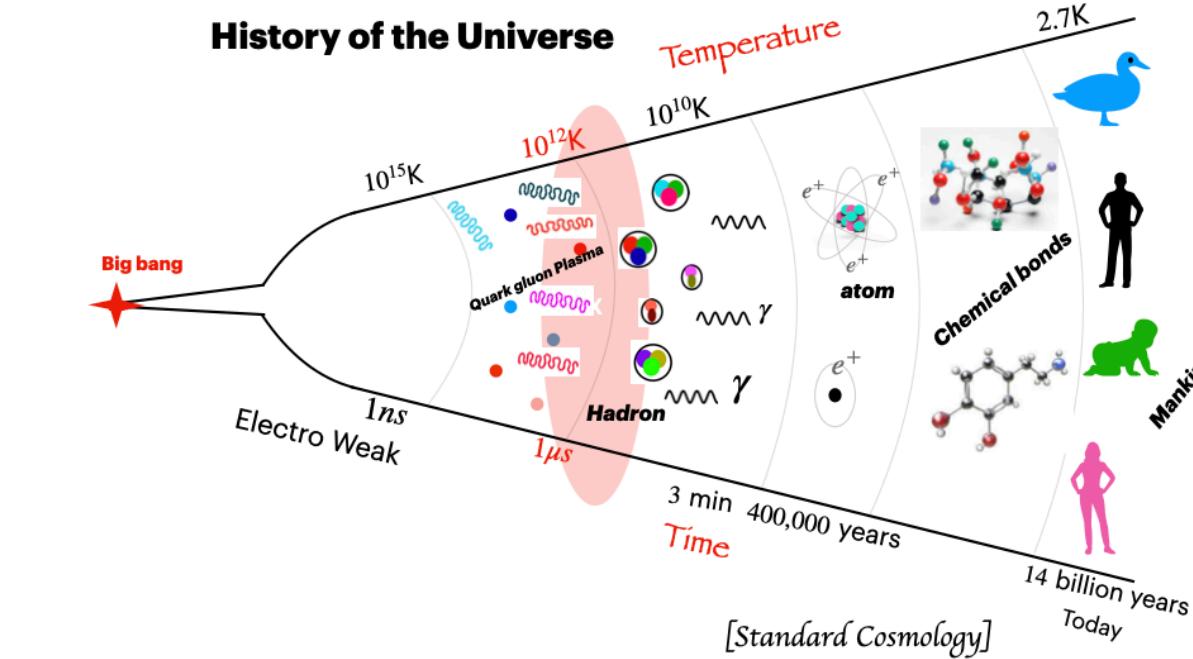
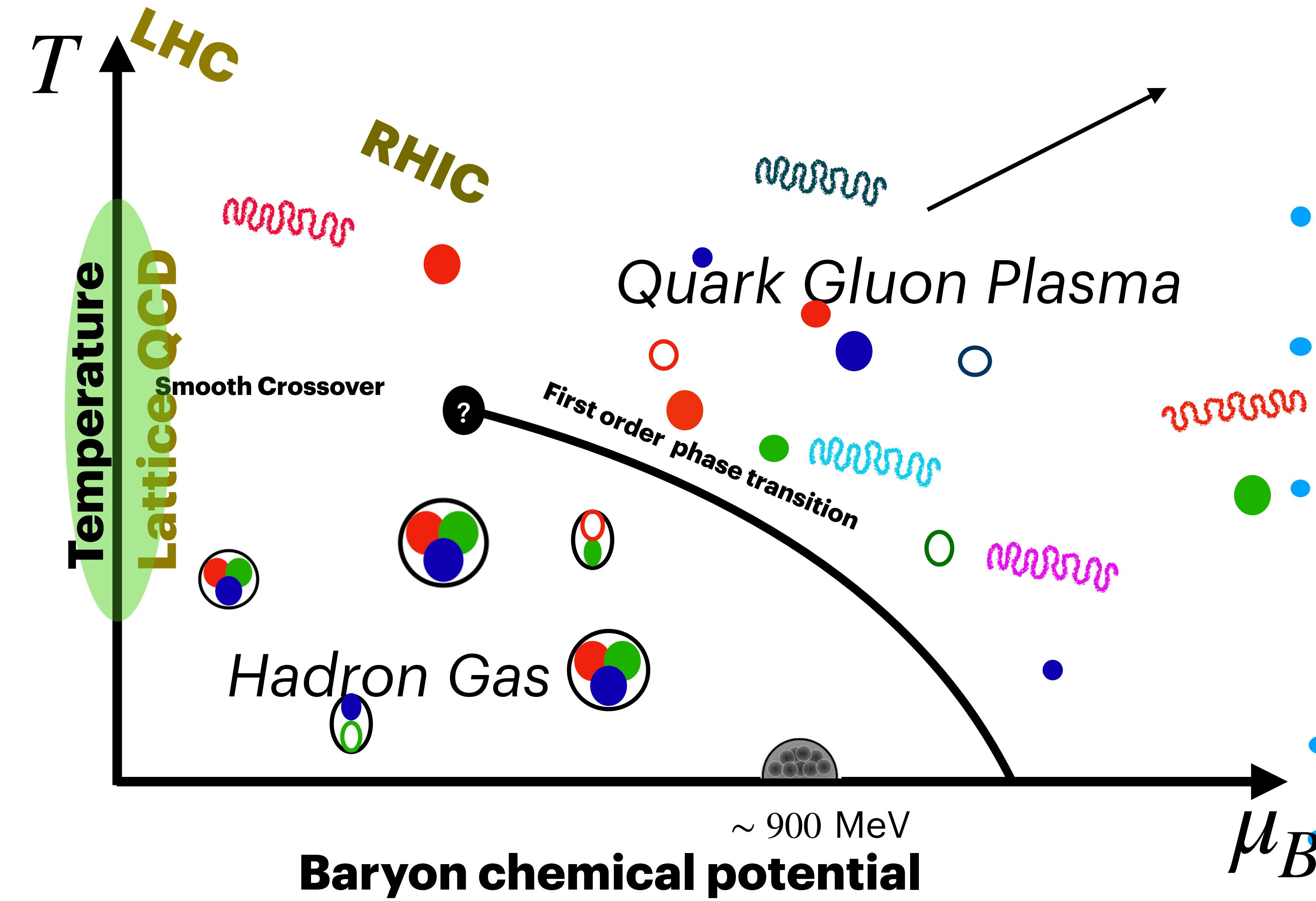


History of the Universe



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What we know

- At $\mu_B = 0$, de-confinement transition is well established (**Smooth crossover**) [[Aoki et al Nature. \(2006\)](#)]
- At finite μ_B , QCD **critical point** is expected but not yet seen
- Lattice simulations are challenging at Finite density (**Fermi-sign problem**)

Attempts

- Experimental programs RHIC,LHC
- Expansion schemes** are used for finite density physics

Taylor: Lattice QCD results

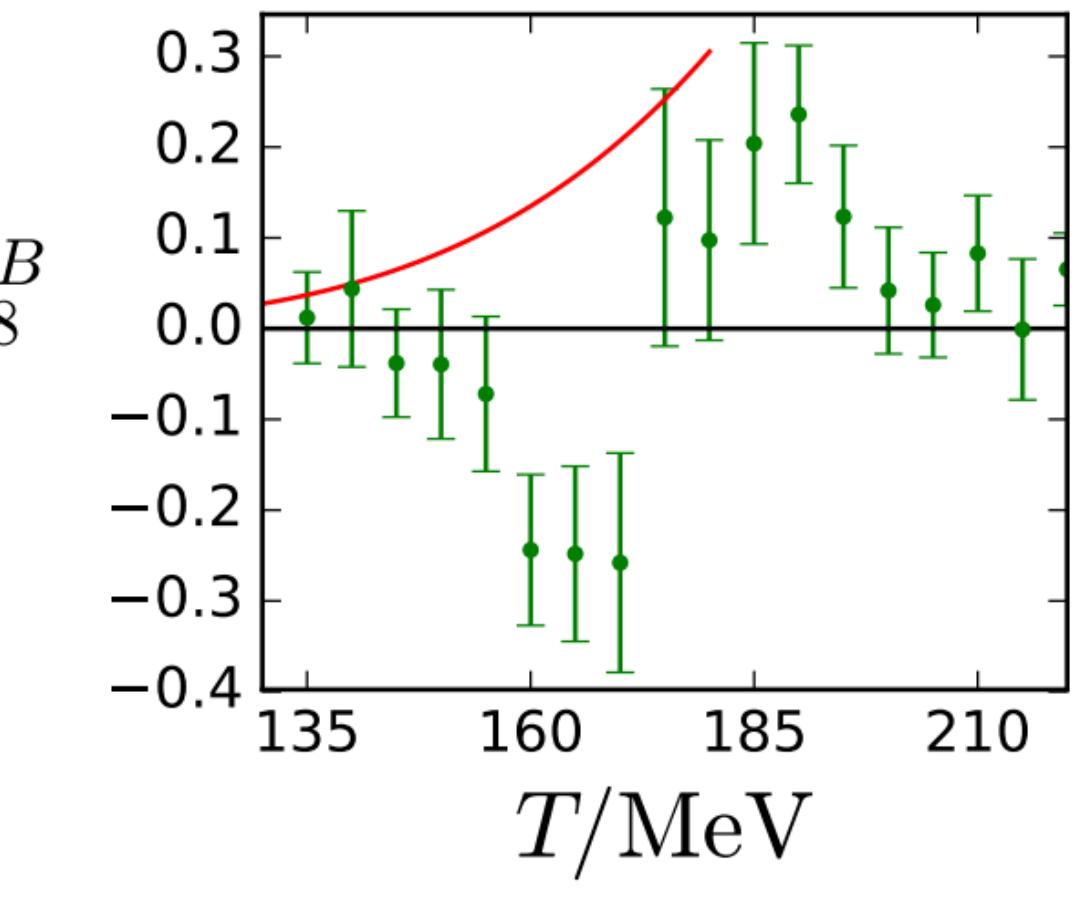
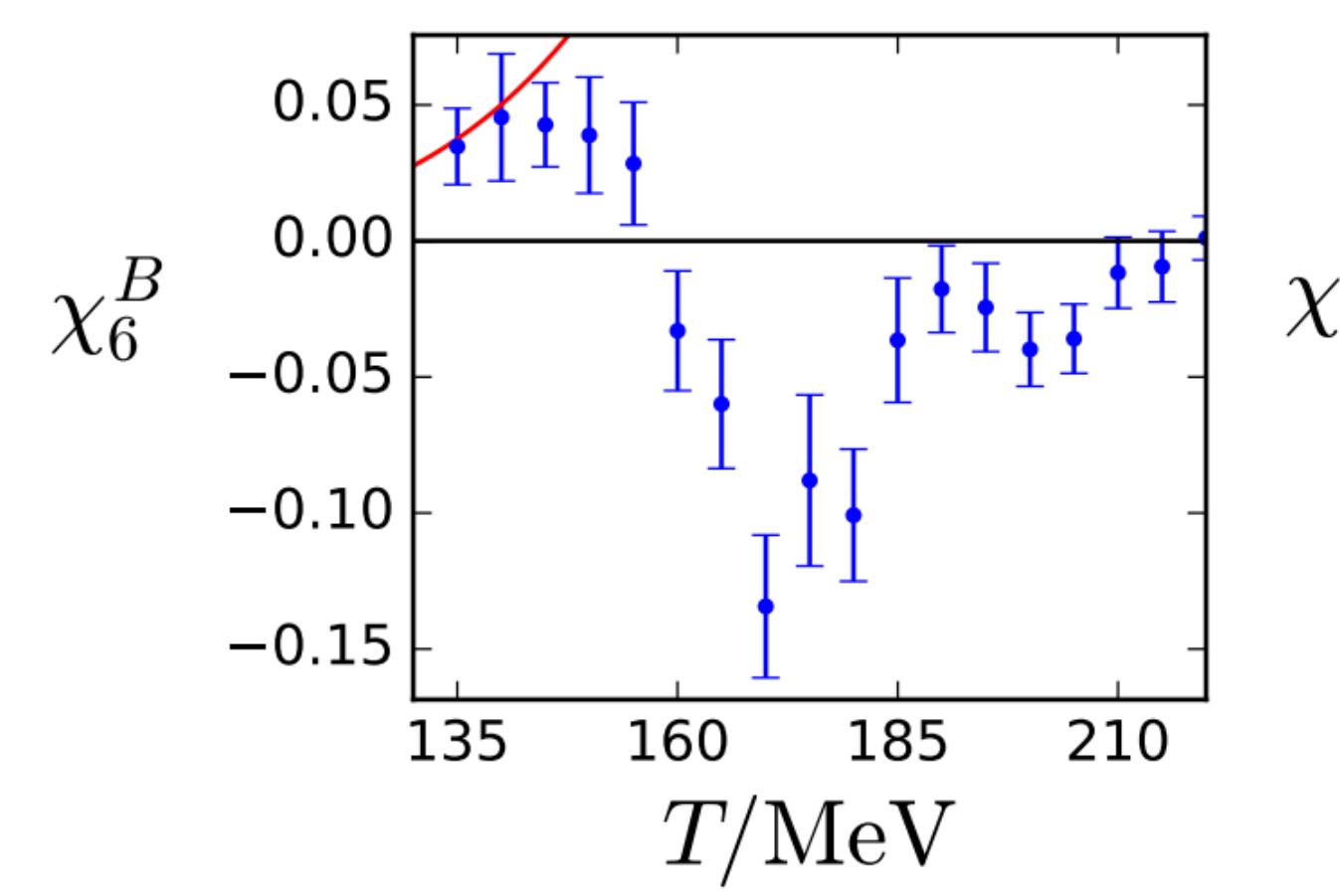
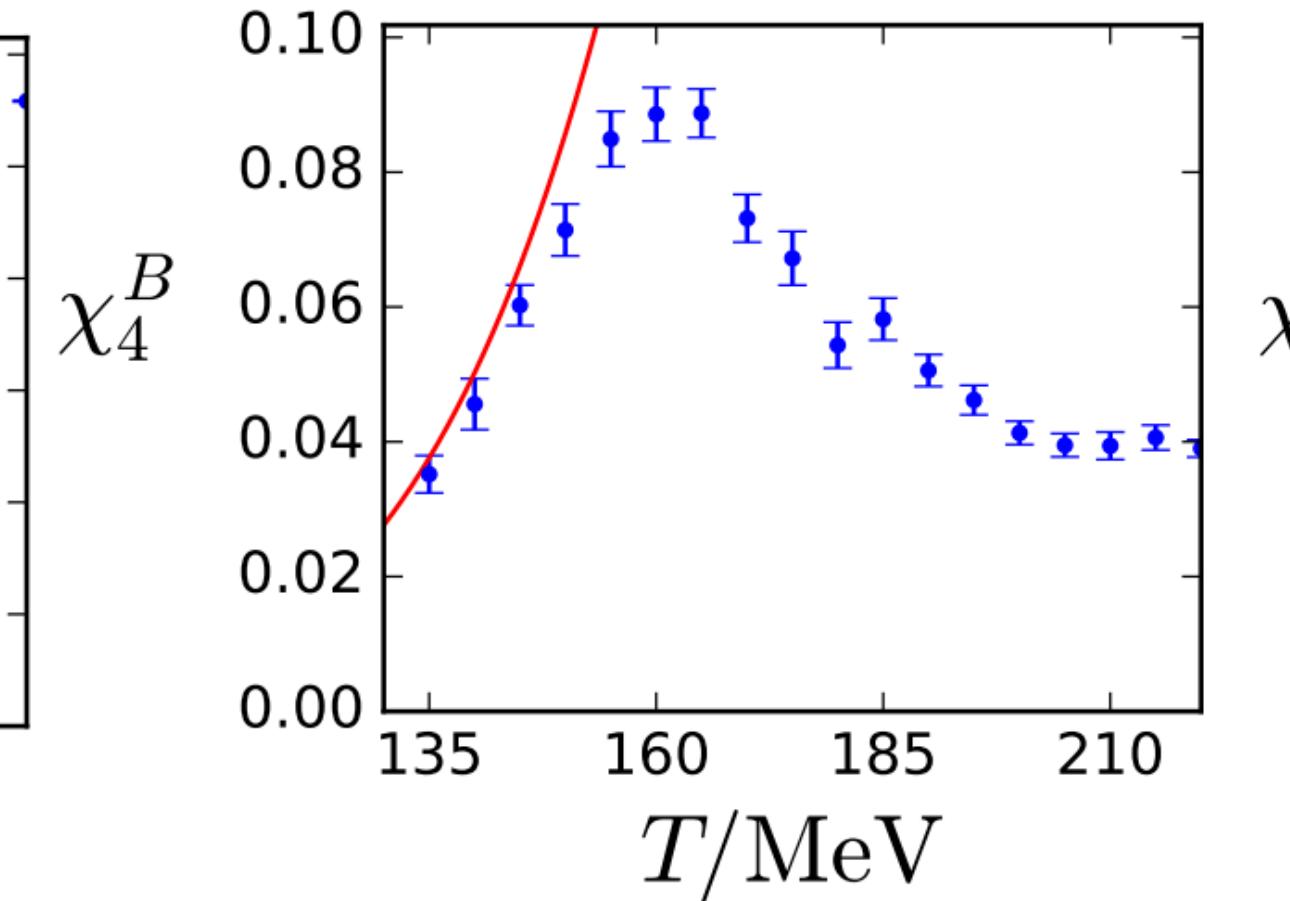
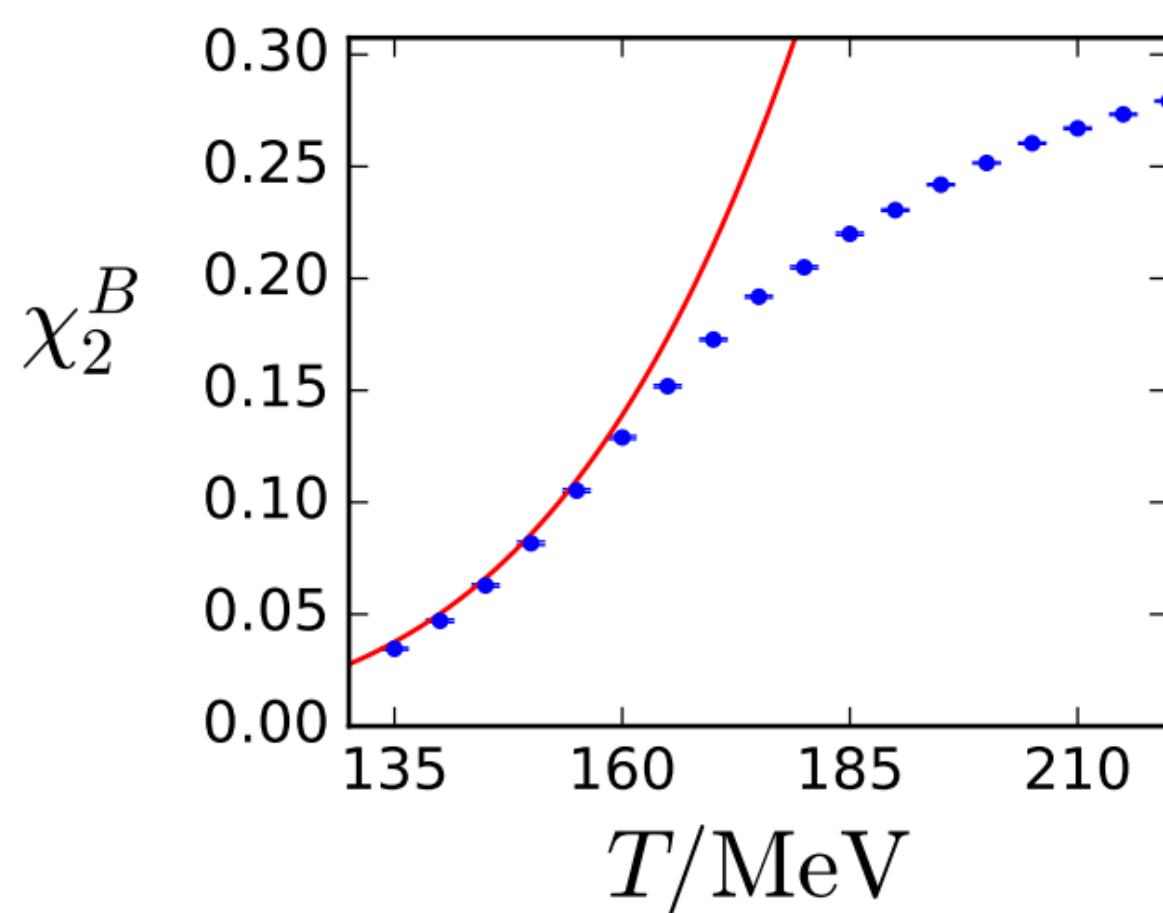
Taylor Expansion around $\mu_B = 0$

$$\frac{n_B(T, \mu_B)}{T^3} = \sum_{n=0}^{\infty} \frac{1}{(2n-1)!} \chi_{2n}(T, \mu_B = 0) \left(\frac{\mu_B}{T} \right)^{2n-1}$$

[Borsanyi, S. et al *High Energy Physics*.9(8), 1-16.(2012)]

[Bazavov, A et al *PhysRevD*.95, 054504 (2017)]

$$c_n(T) = \frac{\chi_n^B(T, \mu_B = 0)}{n!} = \frac{1}{n!} \left(\frac{\partial}{\partial(\mu_B/T)} \right)^n (P/T^4) \Big|_{\mu_B=0}$$



Limitations

- Currently limited to $\frac{\mu_B}{T} \leq 2.5$ despite great computational power
- Including one more higher-order term does not remove unphysical behavior due to truncation of Taylor series

[Bollweg, D. et al *Phys.Rev.D* 108 (2023) 1, 014510]

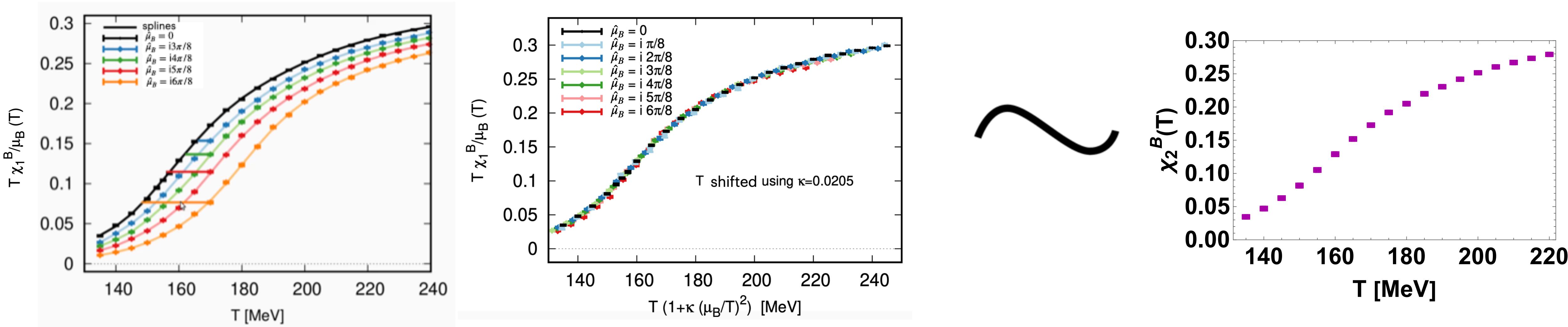
[Borsanyi, S et al *arXiv:2312.07528v1*. (2023)]

[Borsanyi, S. et al *JHEP* 10 205 (2018)]

[Bazavov, A et al *PhysRevD*.95, 054504 (2017)]

T' Expansion scheme (T ExS)

Simulating at Imaginary μ_B



[Borsányi, S et al PRL. 108(1), 101.034901(2021)]

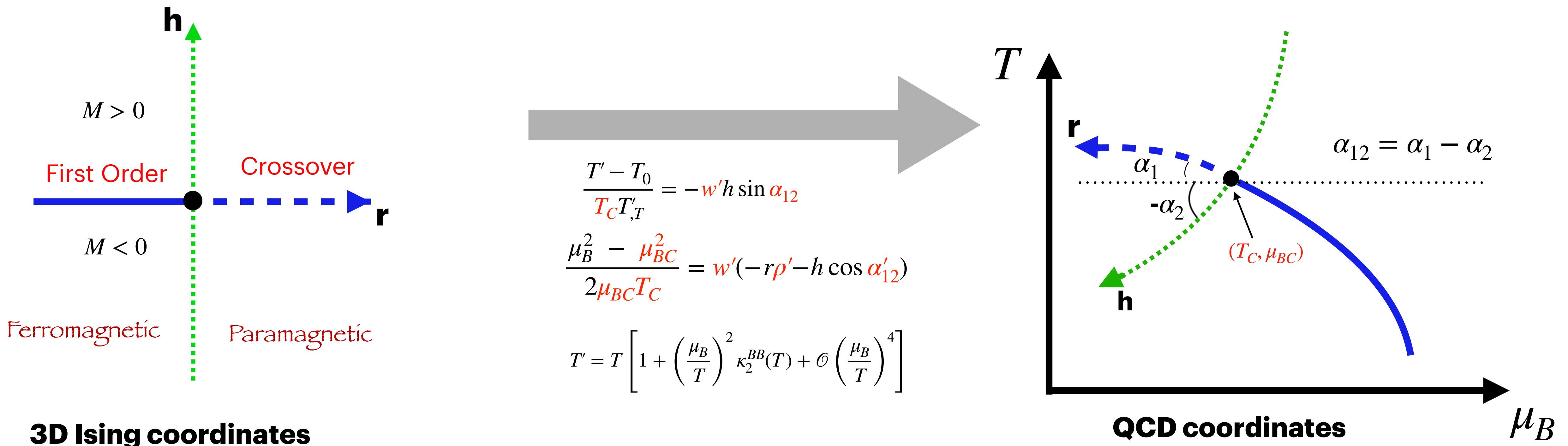
$$T \frac{\chi_1^B(T, \mu_B)}{\mu_B} = \chi_2^B(T, 0)$$

$$T'(T, \mu_B) = T \left[1 + \kappa_2^{BB}(T) \left(\frac{\mu_B}{T} \right)^2 + \kappa_4^{BB}(T) \left(\frac{\mu_B}{T} \right)^4 + \mathcal{O} \left(\frac{\mu_B}{T} \right)^6 \right]$$

- Uses few expansion terms
- μ_B dependence is captured in T-rescaling.
- Trusted up to $\frac{\mu_B}{T} = 3.5$

Introducing Critical Point

Mapping 3D Ising to QCD



[M. K et al arXiv:2402.08636v1]

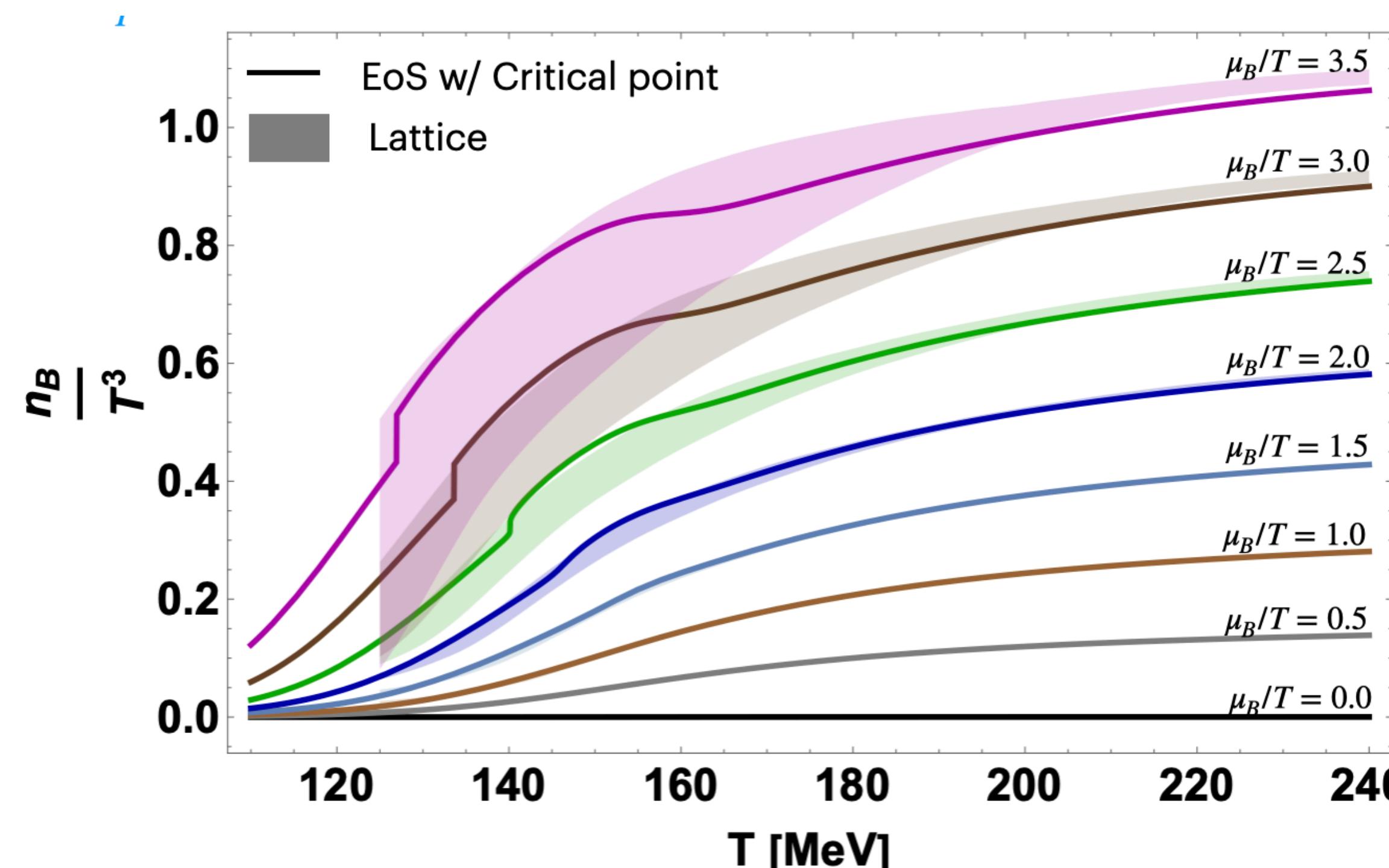
- Free parameters $\mu_{BC}, T_C, w', \rho', \alpha'_{12}$ can be fixed by the current physics knowledge

Merging Ising with Lattice (Ising-T ExS)

Full Baryon Density

$$\frac{n_B(T, \mu_B)}{T^3} = \chi_1^B(T, \mu_B) = \left(\frac{\mu_B}{T} \right) \chi_{2, lat}^B(T', 0)$$

$$T' = T'_{Crit}(T, \mu_B) + T'_{Non-Crit}(T, \mu_B)$$



[M. K et al arXiv:2402.08636v1]

Parameter choice

$$\mu_{BC} = 500 \text{ MeV}$$

$$T_C = 117 \text{ MeV}$$

$$\alpha_{12} = \alpha_1 = 11^0$$

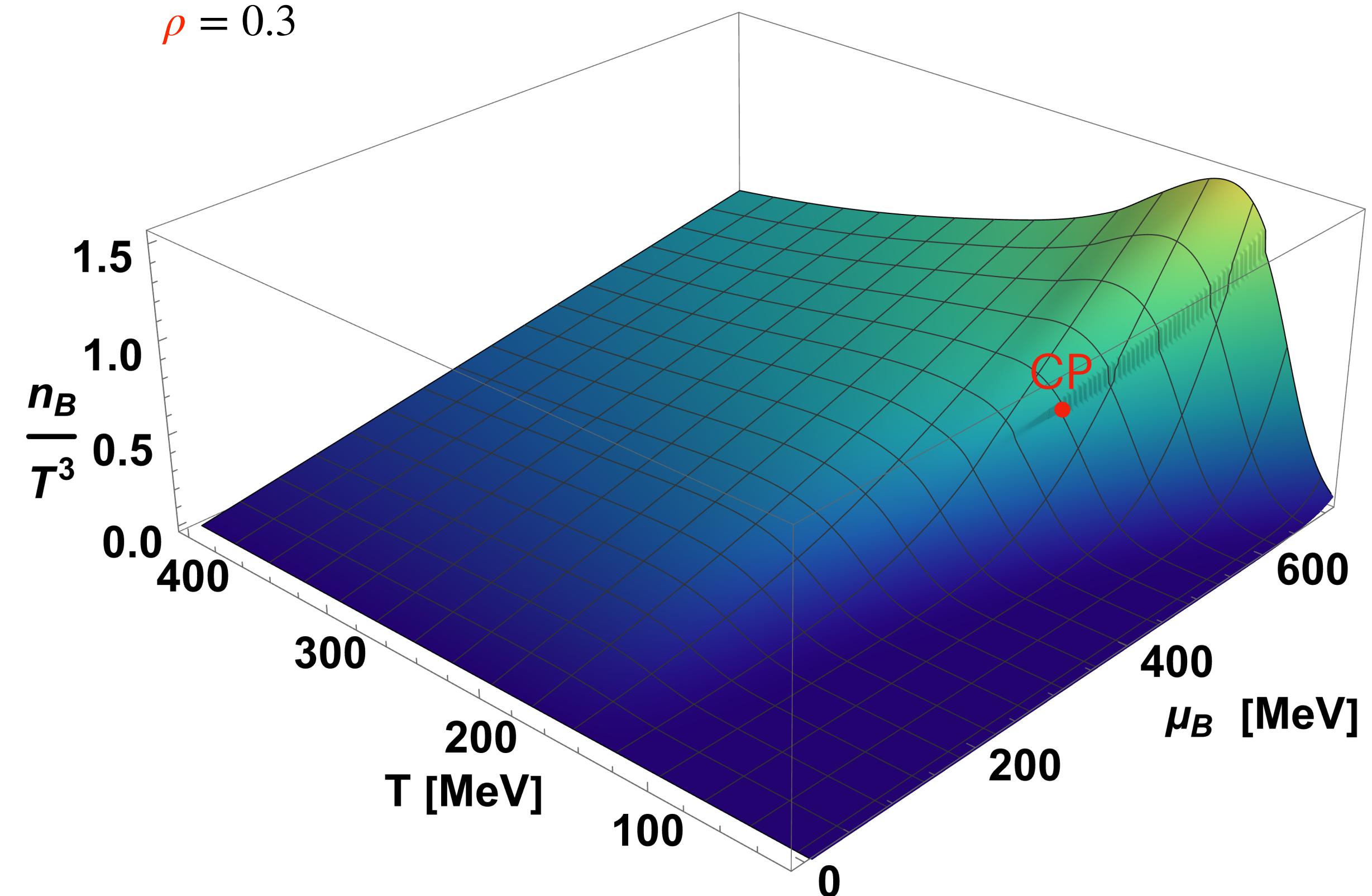
$$\alpha_2 = 0^0$$

$$w = 15$$

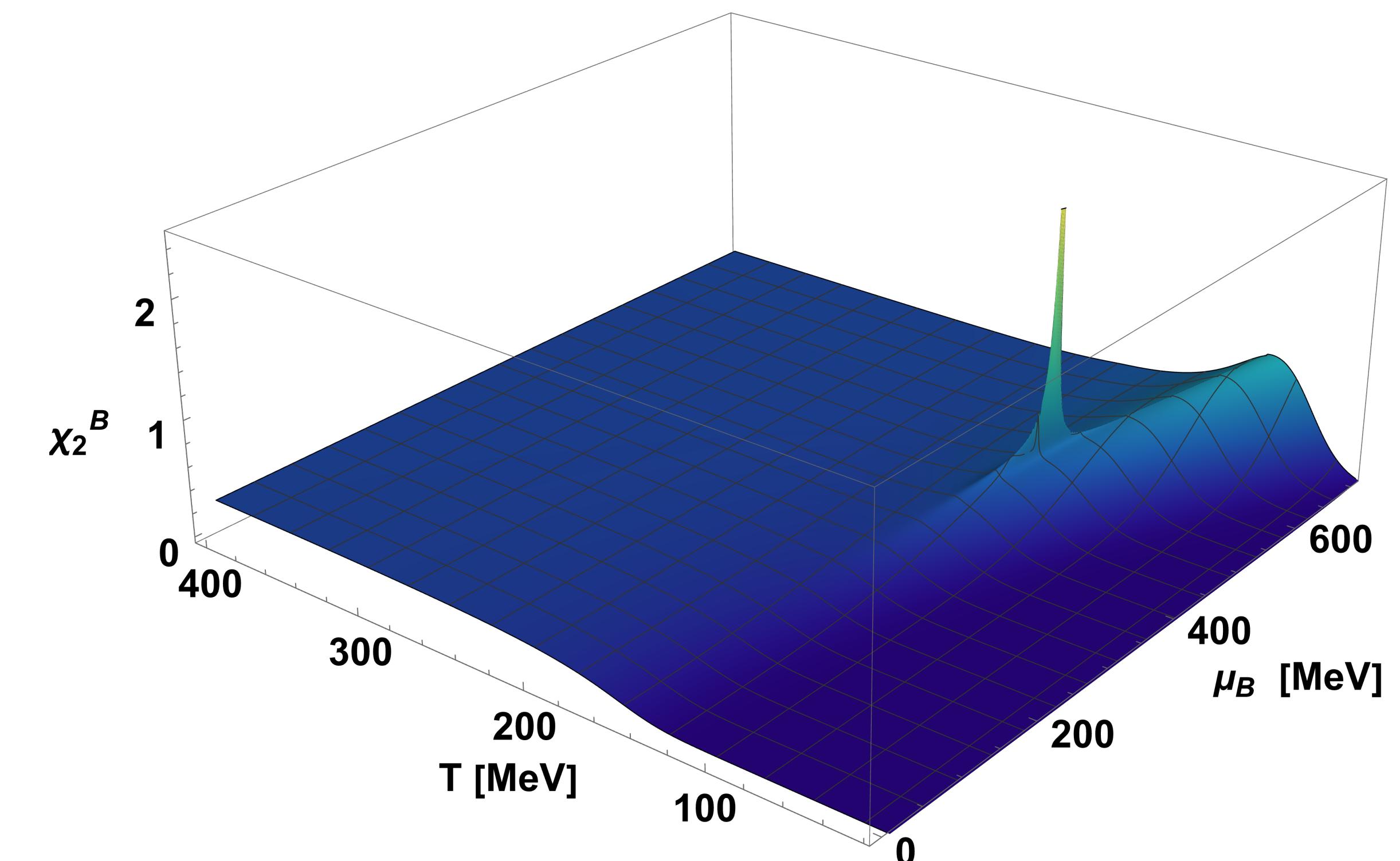
$$\rho = 0.3$$

Other Observables

Baryon Density



Baryon number susceptibility



[M. K et al arXiv:2402.08636v1]

Constraints of the EoS

- Lattice QCD disfavor $\mu_{BC} < 300$ MeV

[Borsányi, S et al PRL. 125, 052001(2020)]

- Choosing μ_{BC} fixes T_C and α_1

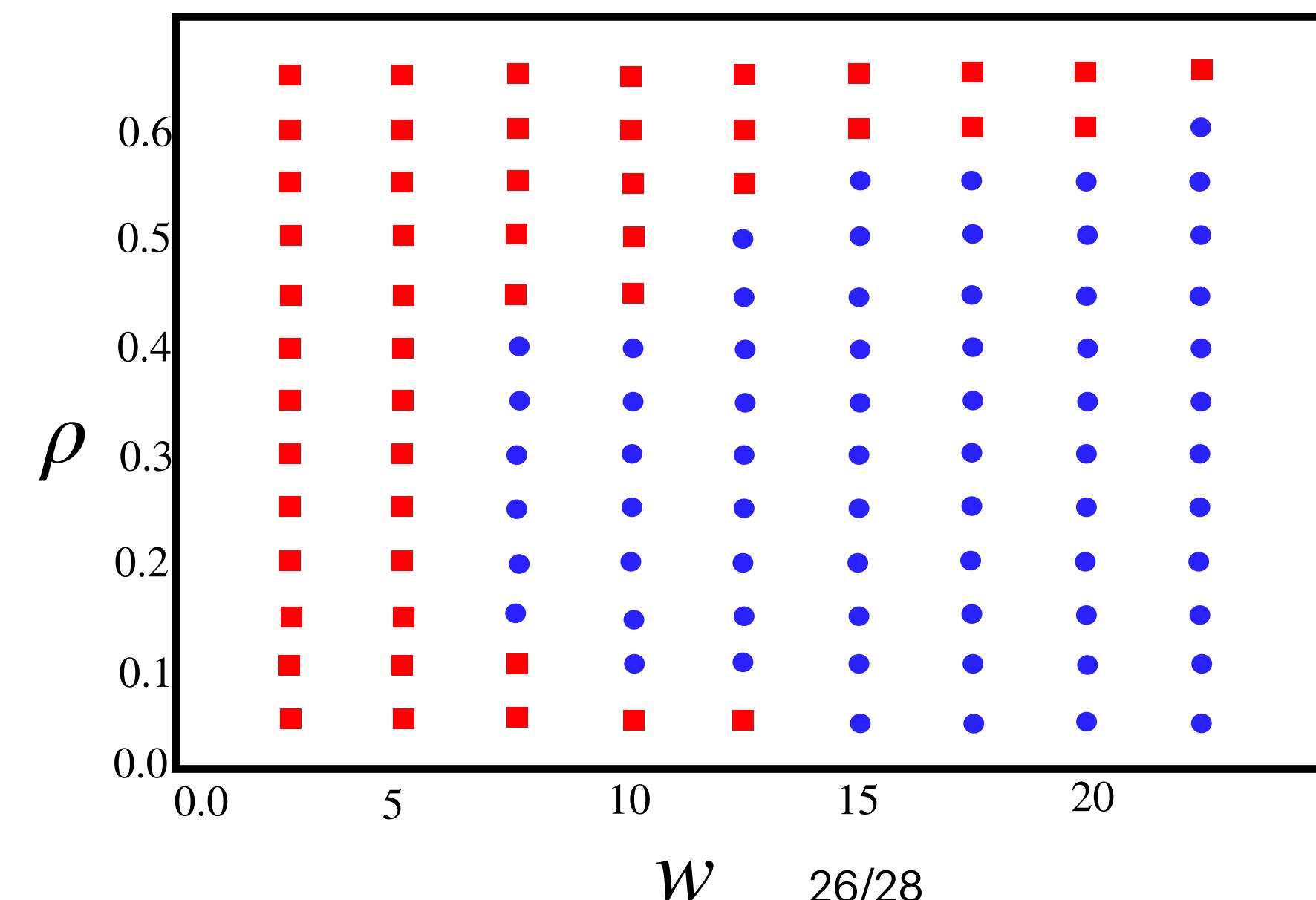
- α_{12} is fixed physical quark masses ($\alpha_{12} = \alpha_1$)

[Pradeep, M. S., & Stephanov, M PhysRevD . 100(5), 056003.(2019)]

- w and ρ are fixed stability and causality

$\mu_{BC} = 500$ MeV

$\alpha_{12} = \alpha_1$



[M. K et al arXiv:2402.08636v1]

Summary and Conclusions

Disclaimer! : We don't predict the location of the critical point

- We provide an Equation of State with enhanced coverage with 3D-Ising model Critical Point



DOI [10.5281/zenodo.10652326](https://doi.org/10.5281/zenodo.10652326)

(Open Software)

Our equation of state, has adjustable parameters, and can be used as input in **hydrodynamical simulations** to compare with experimental searches for the **critical point** in Beam Energy Scan II