QCD equation of state at finite density with a critical point from an alternative expansion scheme

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muses

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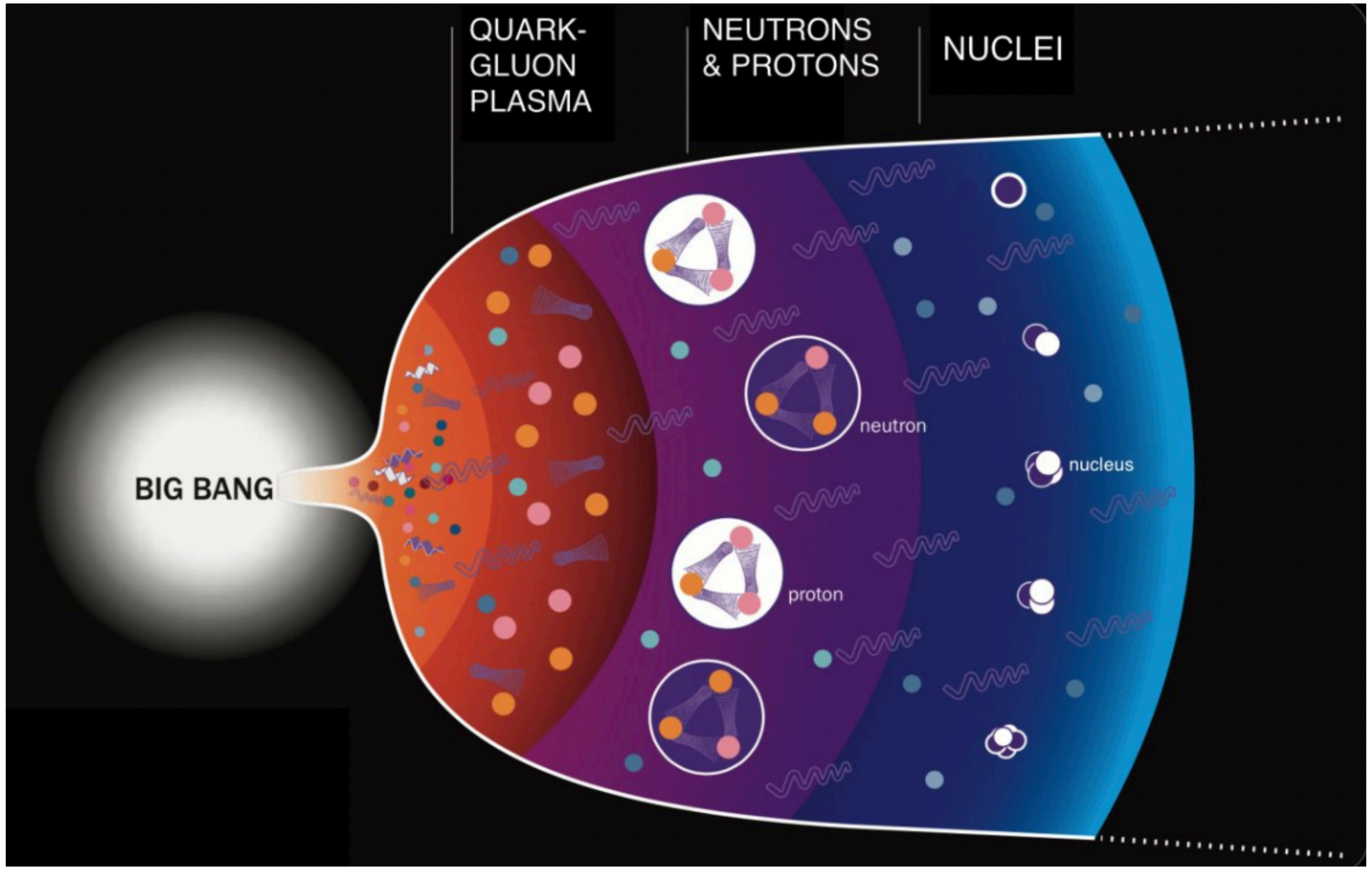
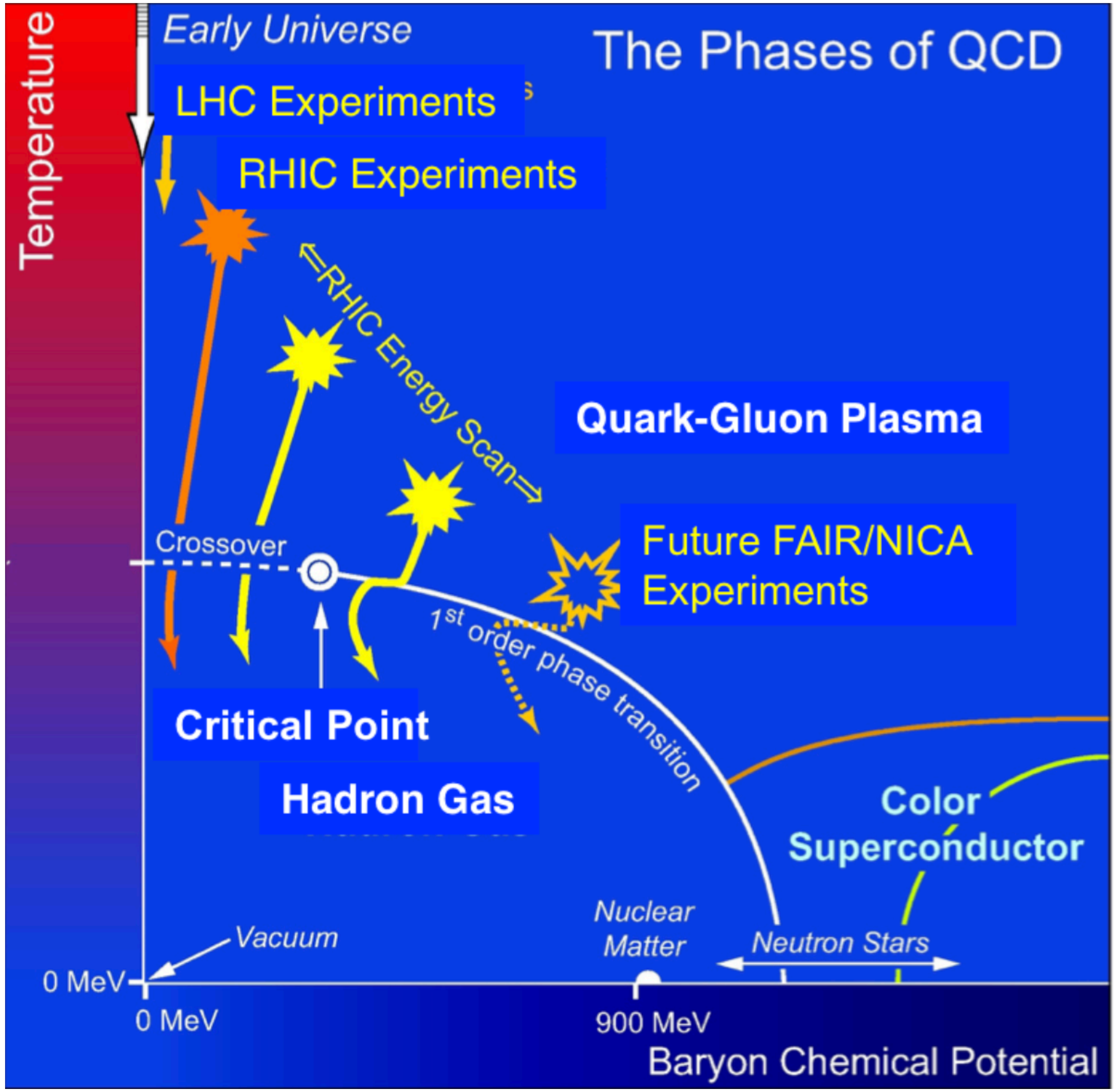
Introduction

QCD Phase Diagram

Current Conjecture





Lattice QCD results

Fermi Sign Problem!

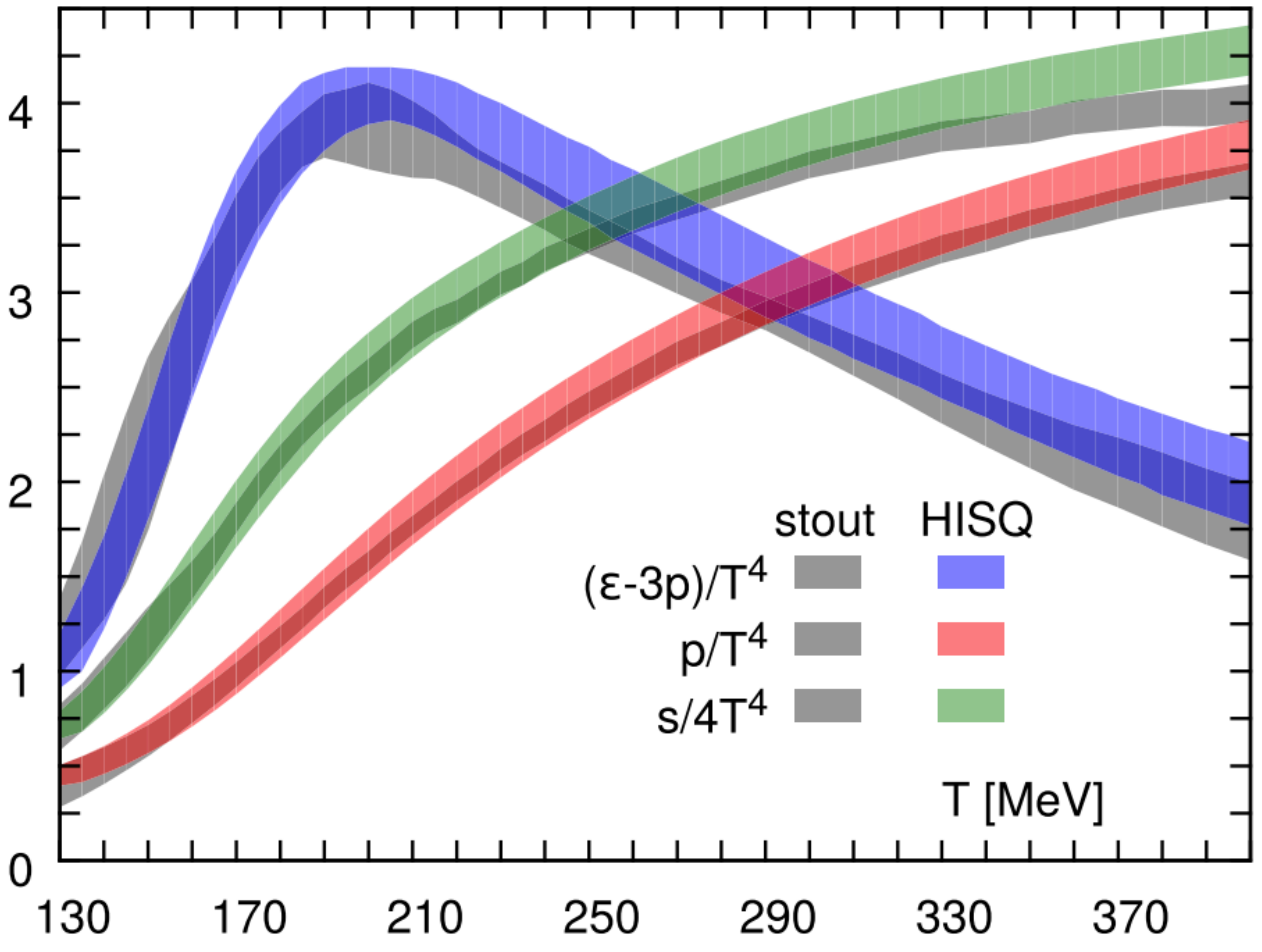
- **Simulation at finite μ_B is Impossible**

$\mu_B = 0$

Way out!

- **Simulating Imaginary μ_B $\rightarrow$ Analytical Continuation**
- **Expansion around $\mu_B = 0$**

[Wuppertal-Budapest (stout) and HotQCD (HISQ) collaborations]



[S. Borsanyi et al Phys. Lett. B 730 (2014) 99]

[A. Bazavov et al.,PhysRevD.90.094503(2014)]

Part 1: Lattice EoS: Taylor Expansion

Lattice EoS at Finite Density

Taylor Expansion around $\mu_B = 0$

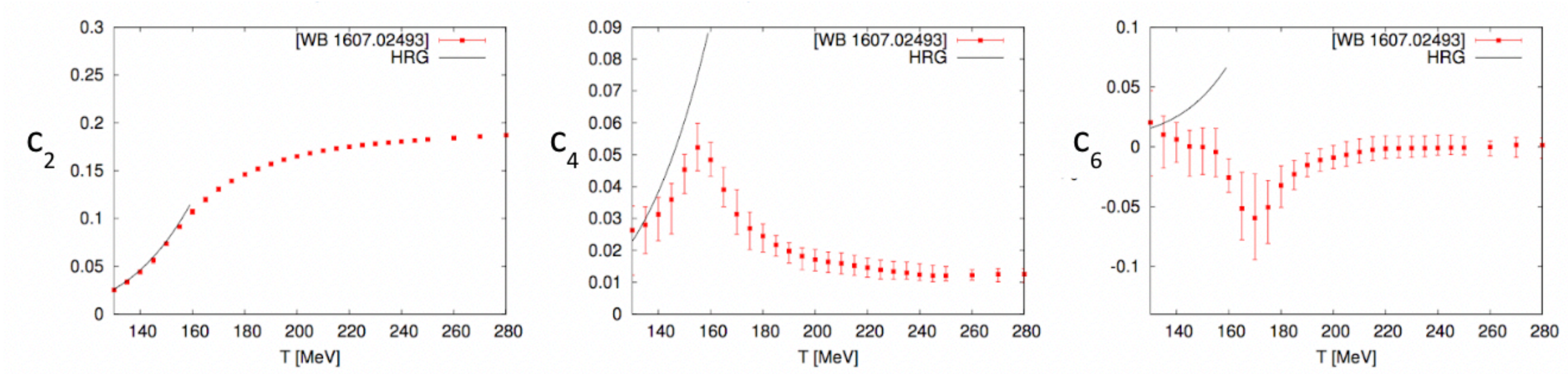
Most Straightforward

$$\frac{P(T,\mu_B)}{T^4} = \sum_{n=0}^{\infty} \frac{1}{n!}\chi_n(T,\mu_B=0)\left(\frac{\mu_B}{T}\right)^n$$

$$c_n(T) = \frac{\chi_n^B(T,\mu_B)}{n!} = \frac{1}{n!}\left(\frac{\partial}{\partial(\mu_B/T)}\right)^n (P/T^4)$$

Limitations

- **Currently limited to $\frac{\mu_B}{T} \leq 2.5$ despite great Computational Power**
- **Adding one more Higher-Order term does not help in convergence**
- **Taylor expansion is carried out at T= constant and doesn't cope well with μ_B-dependent transition temperature**

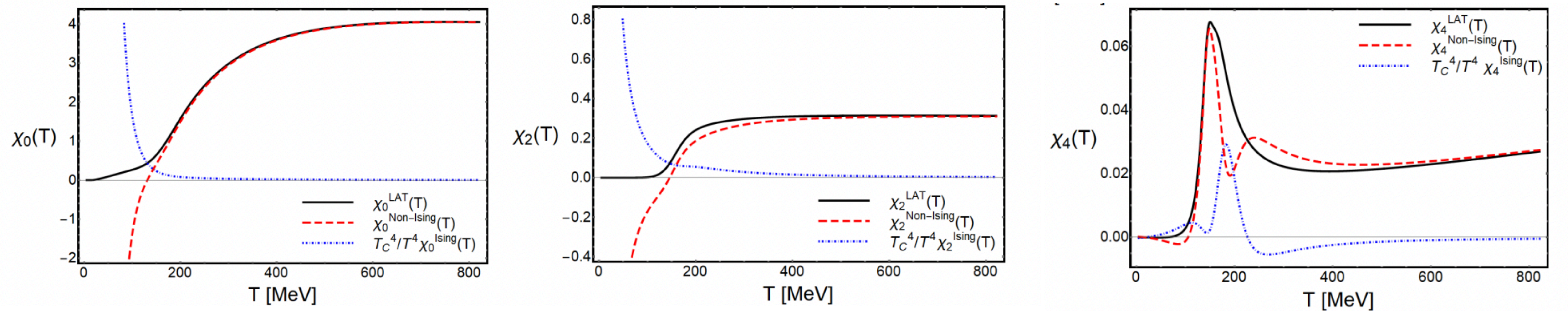


Wuppertal-Budapest

Taylor EoS with Critical point

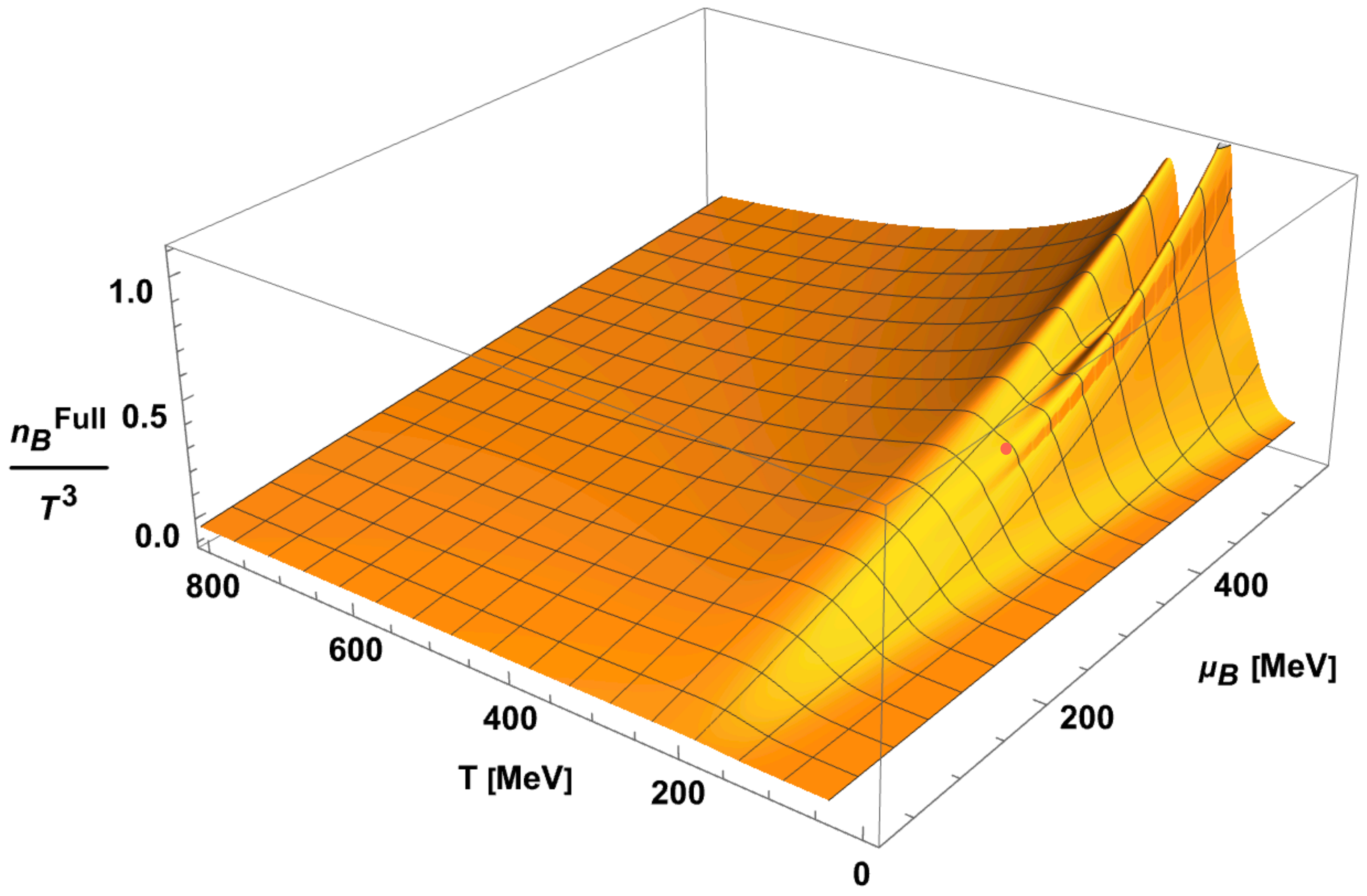
$$P(T,\mu_B) = T^4 \sum_{n=0}^{2} \frac{1}{(2n)!} \chi_{2n}^{Non-Ising}(T) \left(\frac{\mu_B}{T}\right)^{2n} + T_C^4 P_{symm}^{Ising}(T,\mu_B)$$

$$\chi_n^{Lat}(T) = \chi_n^{Non-Ising}(T) + \frac{T_C^4}{T^4} \chi_n^{Ising}(T)$$



3D Diagram

Baryon density



$$\frac{n_B(T,\mu_B)}{T^3} = \frac{\partial(P/T^4)}{\partial(\mu_B/T)}$$

Critical Point at $\mu_{BC} = 350\ [MeV]$

- **Thermodynamic Observables at** $\mu_B \geq 450\ MeV$ **show unphysical behavior**

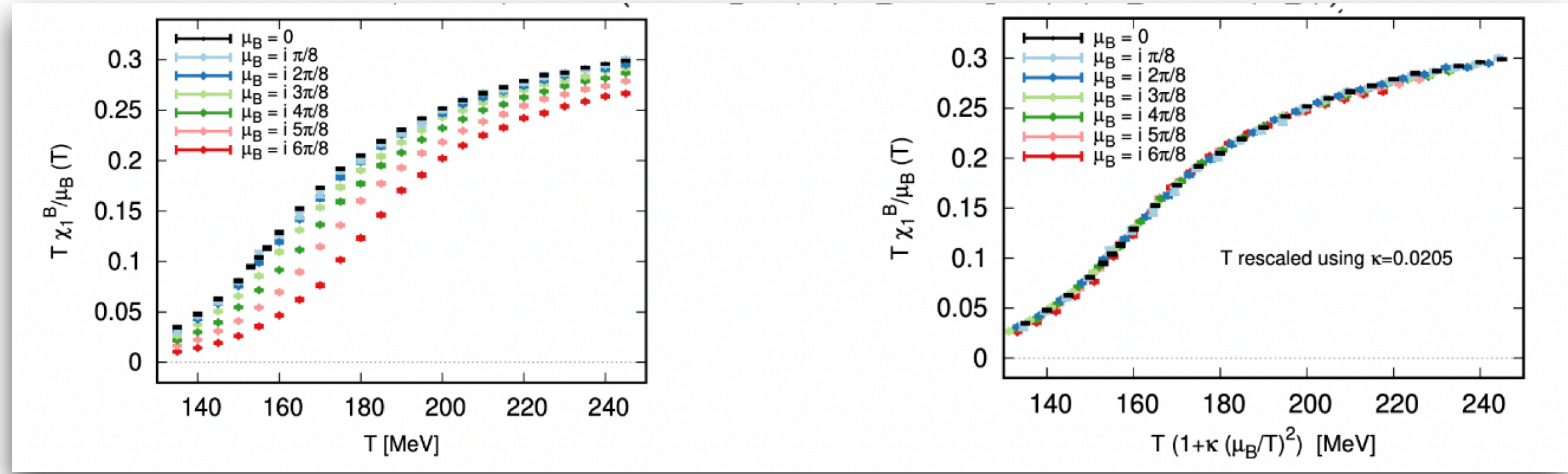
[Paolo Parotto et al PhysRevC.101.034901(2020)]

[BEST Collaboration]

Part 2: Lattice EoS: Alternative Expansion Scheme

Alternative Expansion Scheme: EoS

Simulating at Imaginary μ_B



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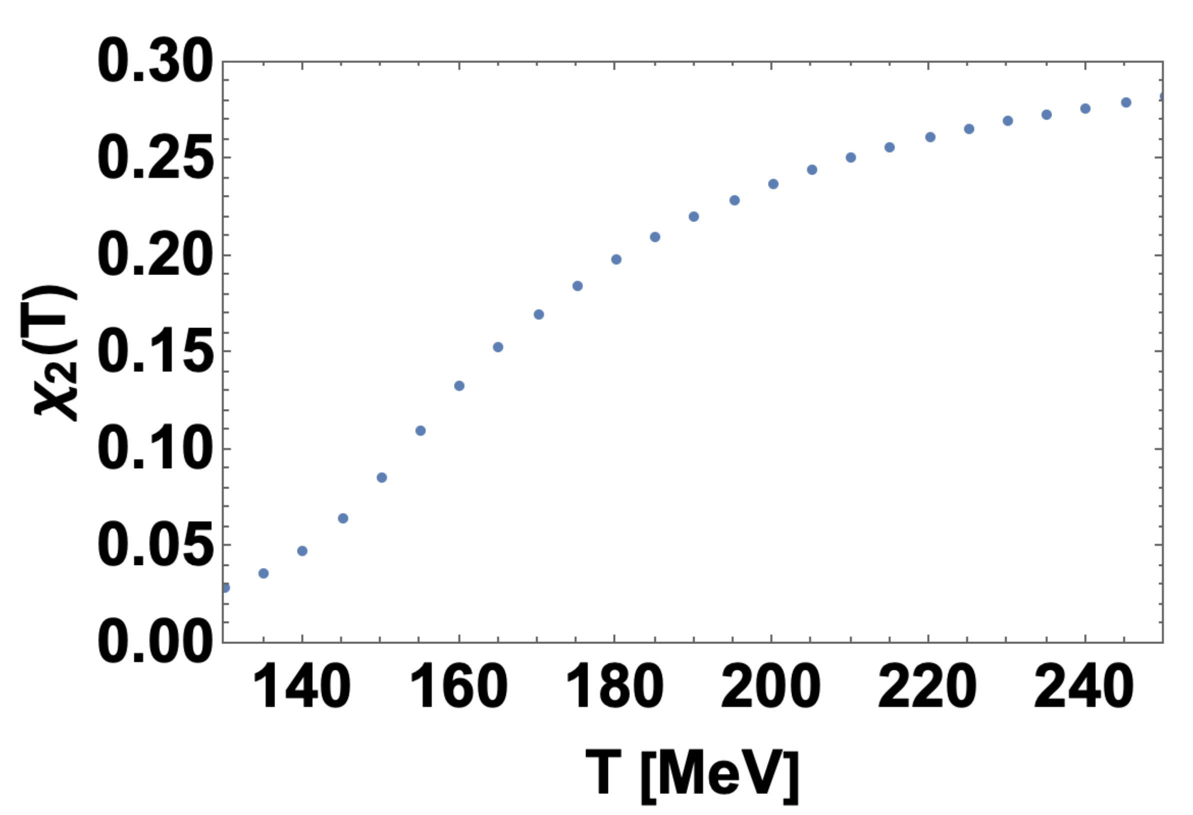


[Borsányi, S. et al. PRL (2021)]

$$T\frac{\chi_1^B(T,\mu_B)}{\mu_B} = \chi_2^B(T',0)$$

$$\chi_n^B(T,\mu_B) = \frac{1}{n!}\left(\frac{\partial}{\partial(\mu_B/T)}\right)^n (P/T^4)$$

$$T'[T,\mu_B] = T\left[1 + \kappa_2^{BB}(T)\left(\frac{\mu_B}{T}\right)^2 + \kappa_4^{BB}(T)\left(\frac{\mu_B}{T}\right)^4 + \mathcal{O}\left(\frac{\mu_B}{T}\right)^6\right]$$

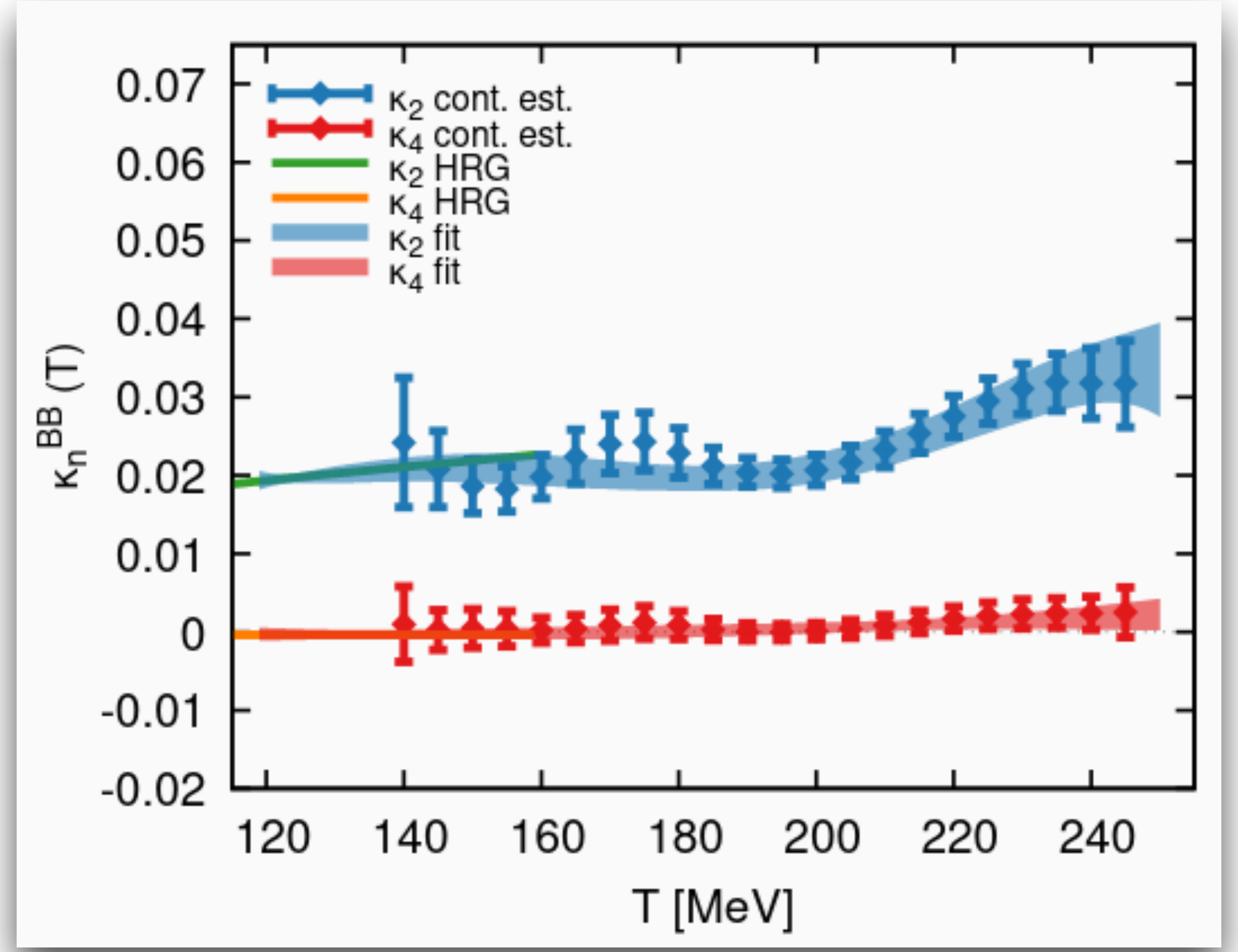
- μ_B **dependence is captured in T-rescaling.**

[Borsányi, S. et al. PRL (2021)]

Alternative expansion scheme

Comparing Taylor expansion and Alternative expansion

- $$\kappa_2^{BB}(T) = \frac{1}{6T}\frac{\chi_4^B(T)}{\left(\frac{\partial \chi_2^B(T)}{\partial T}\right)}$$
- $$\kappa_4^{BB}(T) = \frac{1}{360\chi_2^B(T)^3}\left(3\chi_2'^{B2}\chi_6^B(T) - 5\chi_2^B(T)''\chi_4^B(T)^2\right)$$



[Borsányi, S. et al. PRL (2021)]

Pros

- $\kappa_2(T)$ **is fairly constant over a large T-Range**
- **There is a separation of scale between** $\kappa_2(T)$ **and** $\kappa_4(T)$
- $\kappa_4(T)$ **is almost zero → faster convergence**
- **A good agreement with HRG results at Low Temperature**

Part 3: Putting Critical Point into Alternative Expansion: EoS

EoS with a Critical point

Strategy

Re-scaling EoS for 3D-Ising model → **Map Ising to QCD** → **Estimate the critical Contribution** → **Reconstruct the Full Baryon Density**

Re-scaling EoS for 3D-Ising model

Close to the critical point, we define parametrization for Magnetization M, Magnetic field h, and reduced temperature

QCD Critical point is in the 3D-Ising model Universality class

$$r = \frac{T - T_C}{T_C}$$

$$M = M_0 R^{\beta} \theta$$

$$h = h_0 R^{\beta\delta} \tilde{h}(\theta)$$

$$r = R(1 - \theta^2)$$

$$(R, \theta) \longmapsto (r, h)$$

- M_0, h_0 **are normalization constants**
- $\beta \approx 0.326$ **and** $\delta = 4.8$ **are the 3D Ising model critical exponents**
- $\tilde{h}(\theta) = (\theta + a\theta^3 + b\theta^5)$ **with** $a = -0.76201, b = 0.00804$
- **The parameters take on the values** $R \geq 0, |\theta| \leq \theta_0 \approx 1.154, \theta_0$ **being the first non-trivial zero of** $\tilde{h}(\theta)$

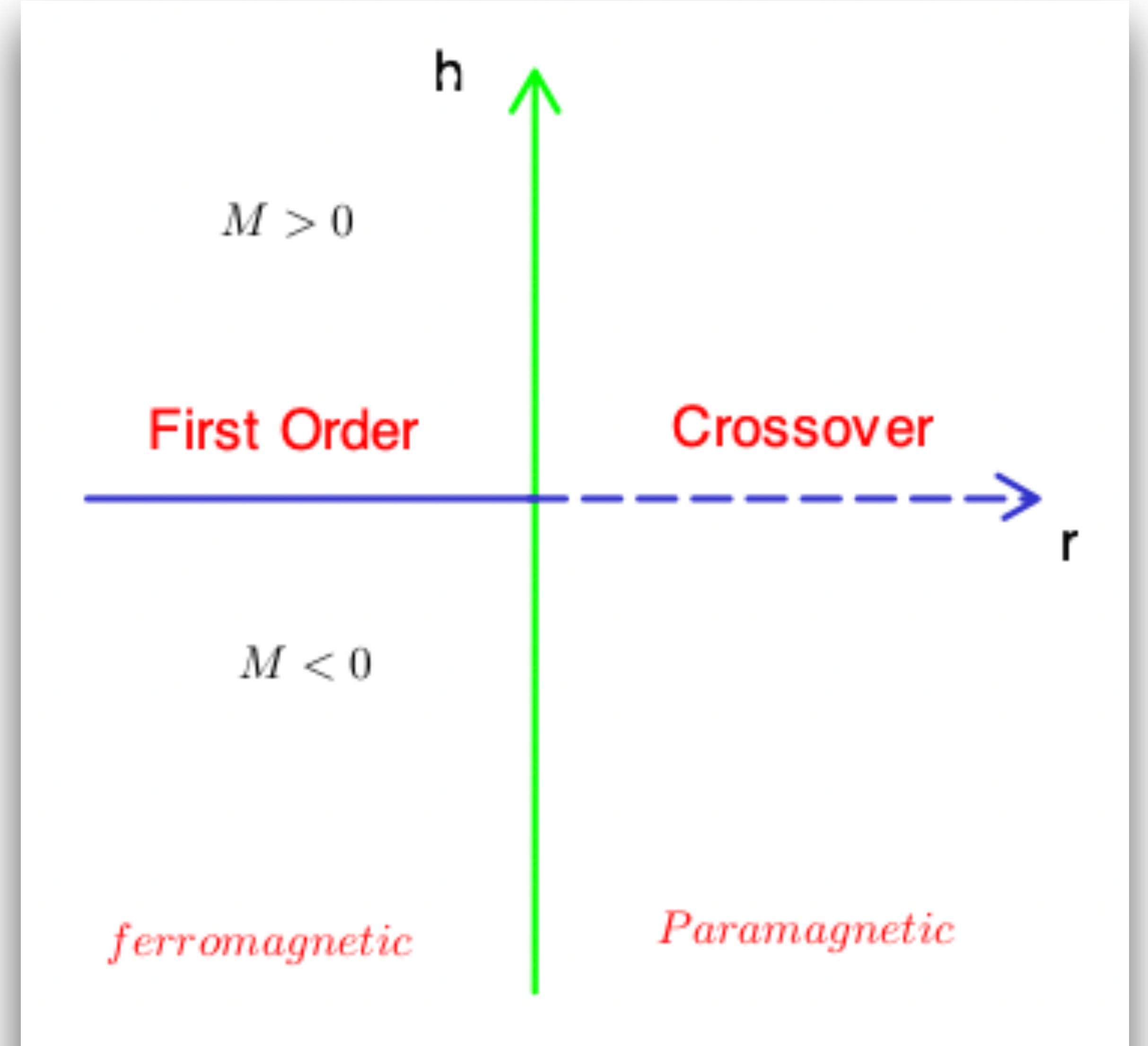
[Parotto et al PhysRevC.101.034901(2020)]

[Nonaka et al Physical Review C, 71(4), 044904.(2005)]

[Guida et al Nuclear Physics B, 489(3), 626-652.(1997)]

EoS with a Critical point

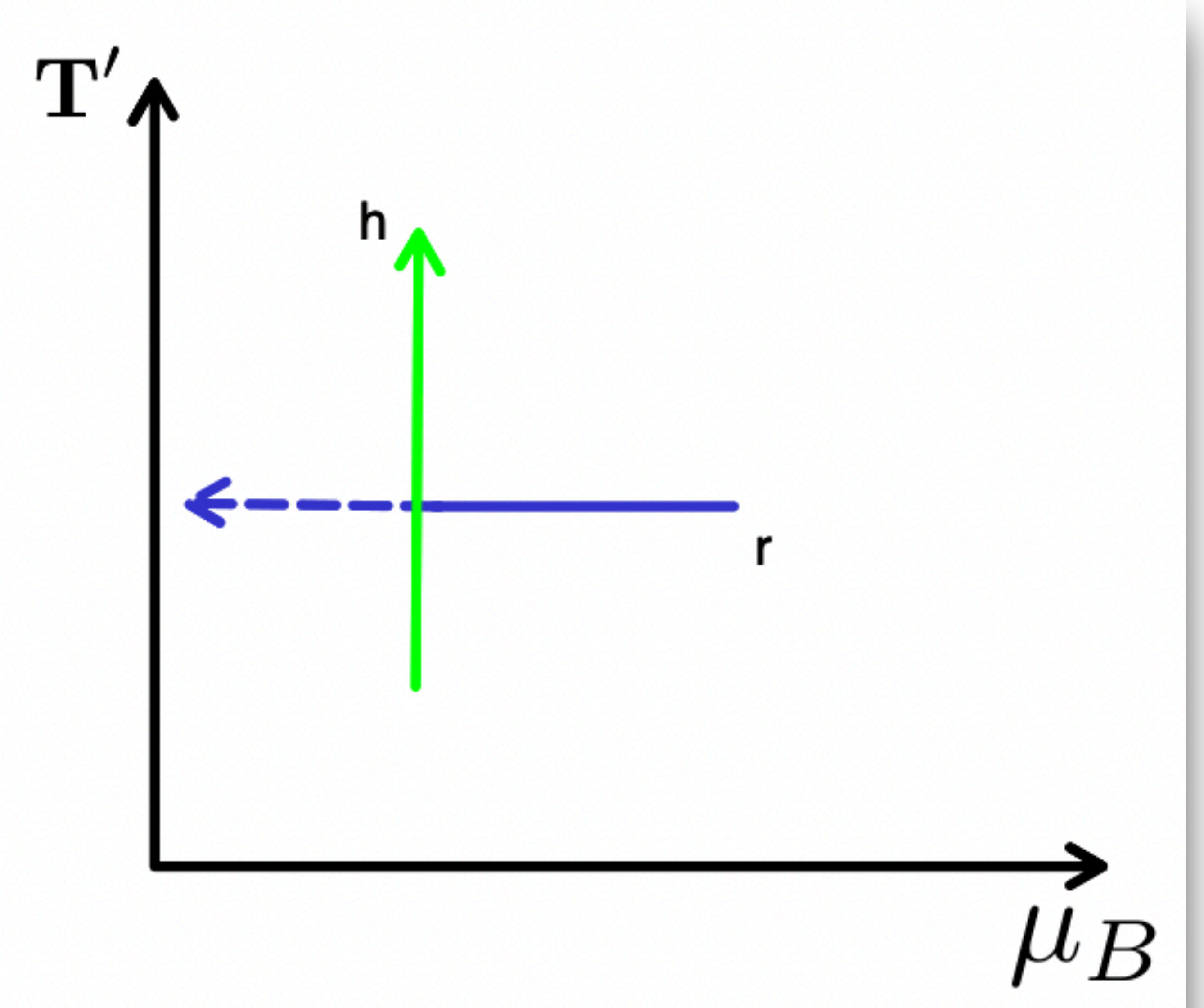
Ising to QCD Mapping



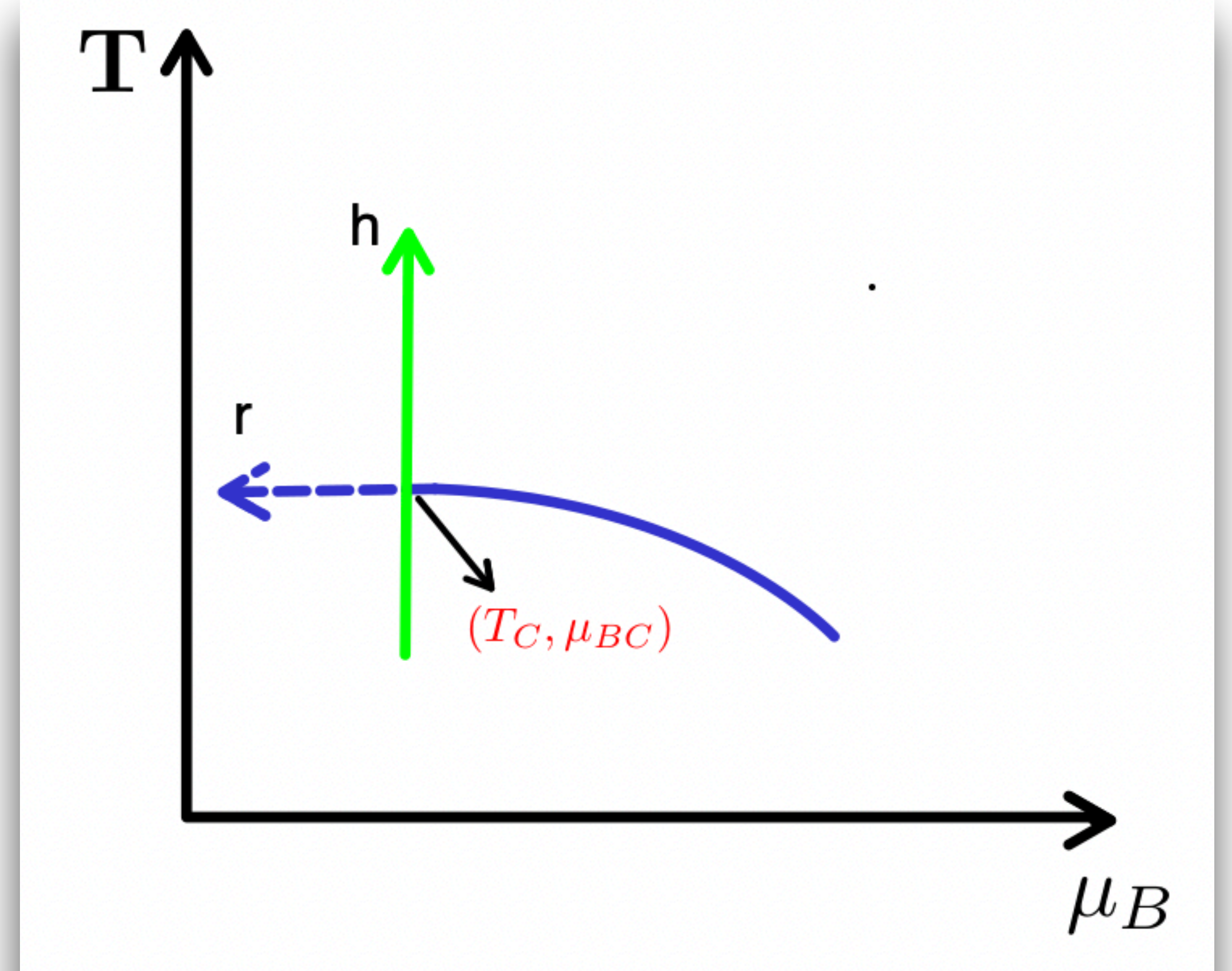
3D-Ising Co-ordinates

$$\frac{T' - T_0}{T_0} = w\ h \sin\ \alpha_{12}$$

$$\frac{\mu_B - \mu_{BC}}{T_0} = w\,(-r\rho - h\cos\ \alpha_{12})$$



(T′, μ_B)-Co-ordinates

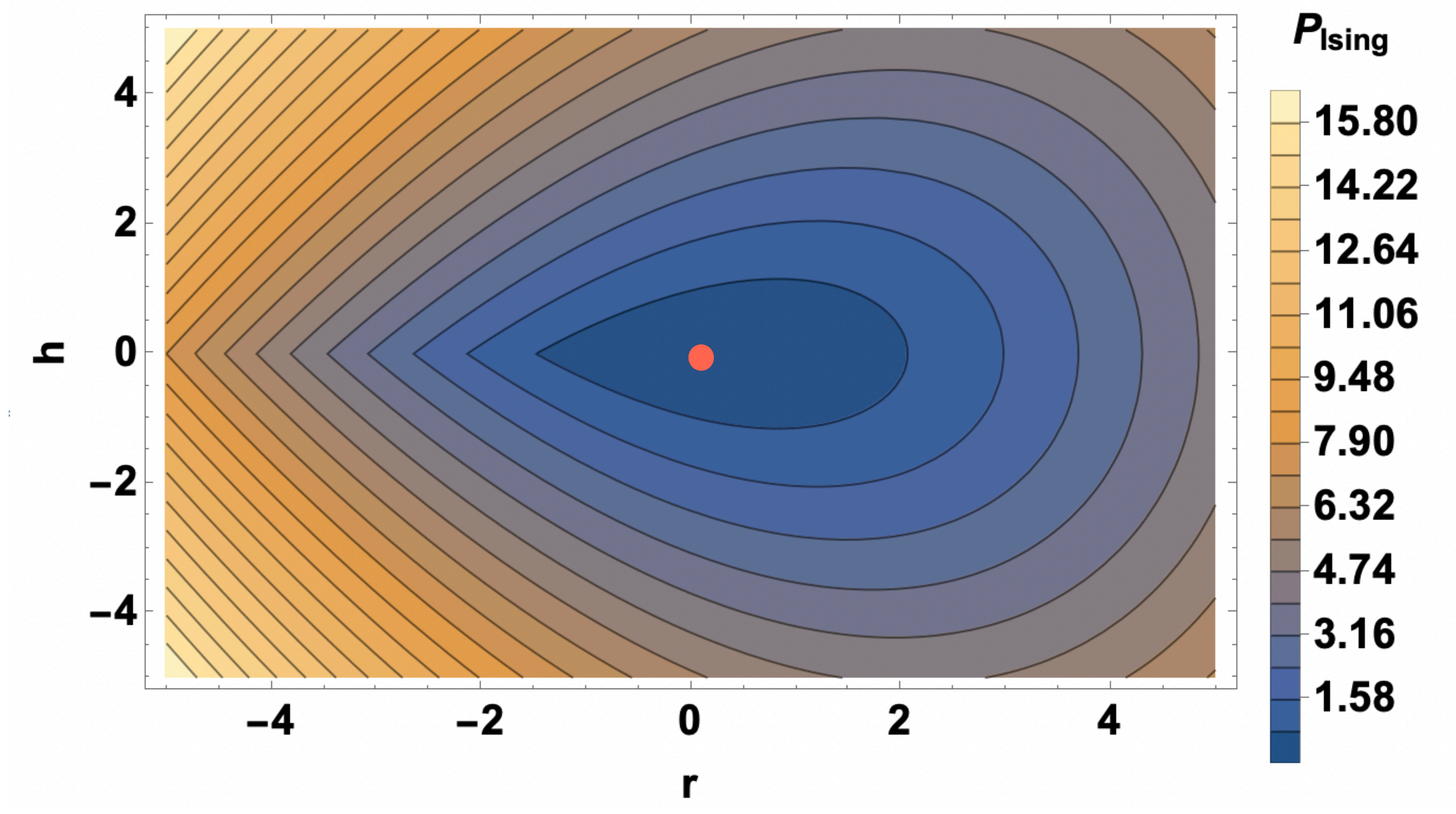


QCD - Co-ordinates

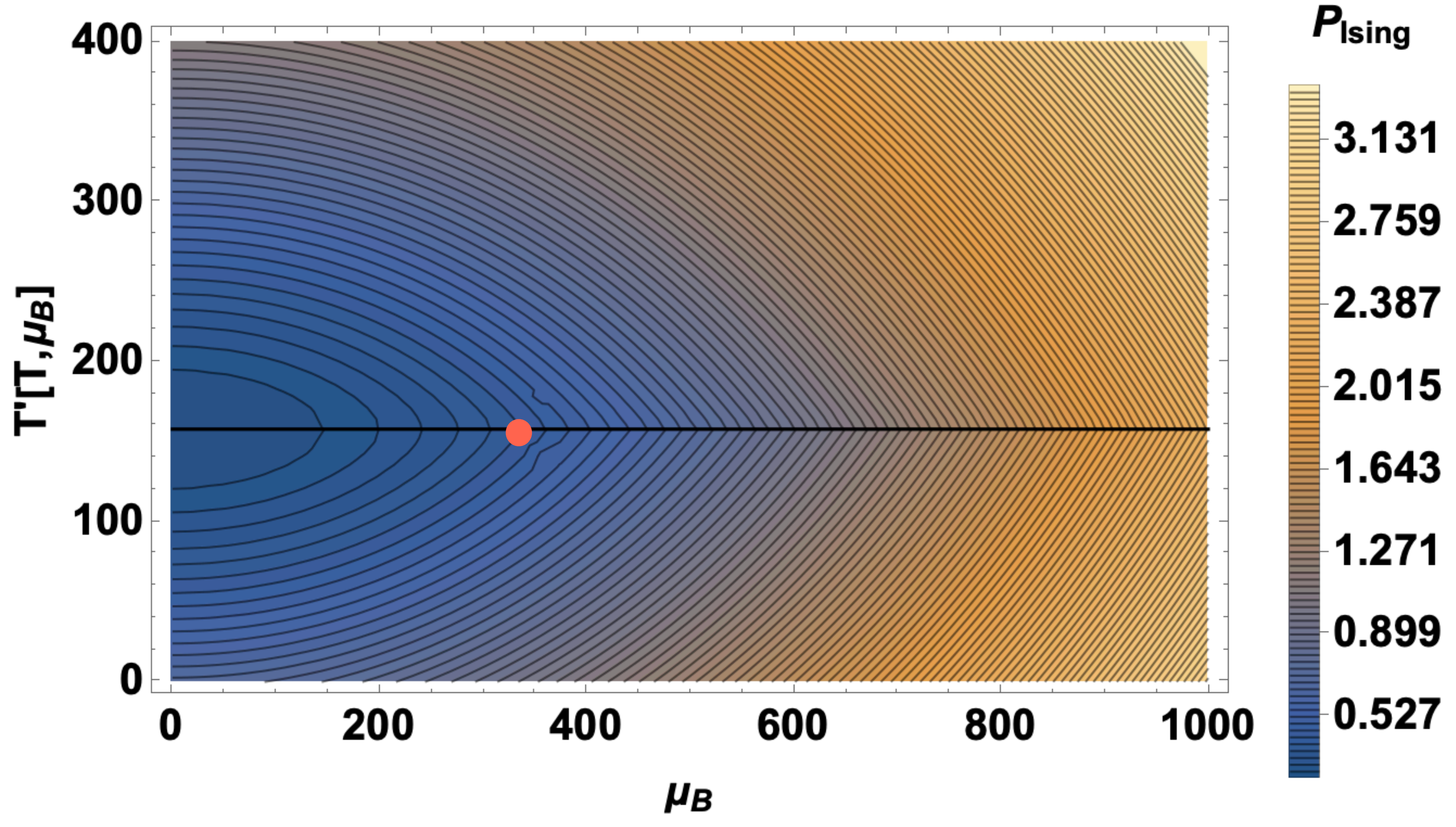
$$T' = T\left[1 + \left(\frac{\mu_B}{T}\right)^2 \kappa_2^{BB}(T) + \mathcal{O}\left(\frac{\mu_B}{T}\right)^4\right]$$

w(-r\rho - h \cos\alpha_{12})

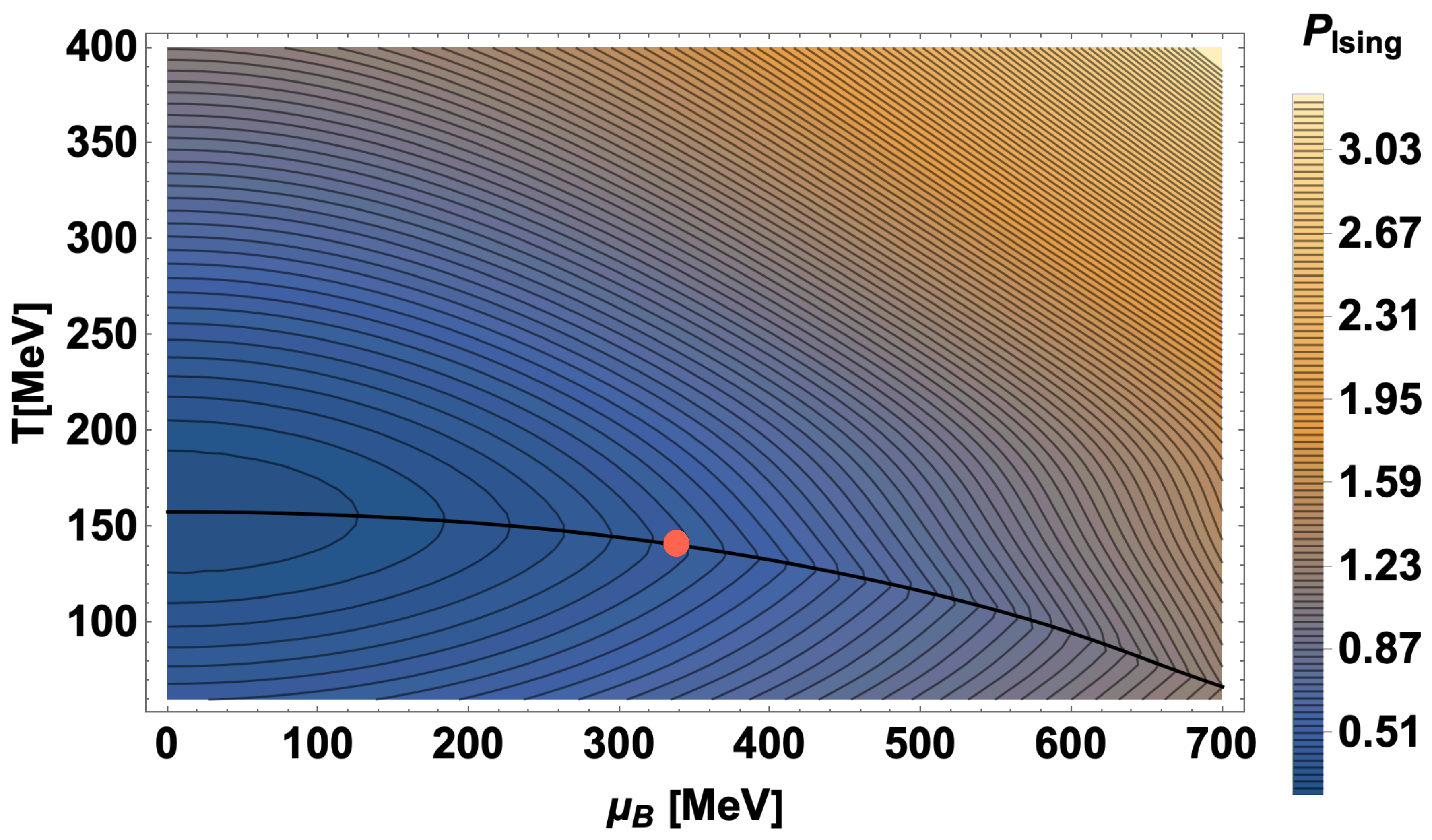
Ising Pressure



Ising Coordinates



(μ_B, T') **Coordinates**



QCD Coordinates

Parameters

$$w = 1,\ \rho = 2,\ \mu_{BC} = 350\ MeV,\ T_0 = 158\ [MeV]$$

$$T_C(\mu_B) = T_0 \left[1 - \kappa_2(T_0) \left(\frac{\mu_B}{T_0} \right)^2 \right]$$

Re-Constructing the Full Baryon Density

$$\frac{n_B^{full}(T,\mu_B)}{T^3} = \left(\frac{\mu_B}{T}\right) \chi^B_{2,Lattice}(T'_{full}(T,\mu_B),0)$$

$$T'_{full}(T,\mu_B) \quad = \quad \underbrace{T'_{Lattice}(T,\mu_B)}_{\text{lowest order in } (\frac{\mu_B}{T})} \quad + \underbrace{T'_{crit}(T,\mu_B) - Taylor[T'_{crit}(T,\mu_B)]}_{\text{higher orders in } (\frac{\mu_B}{T})}$$

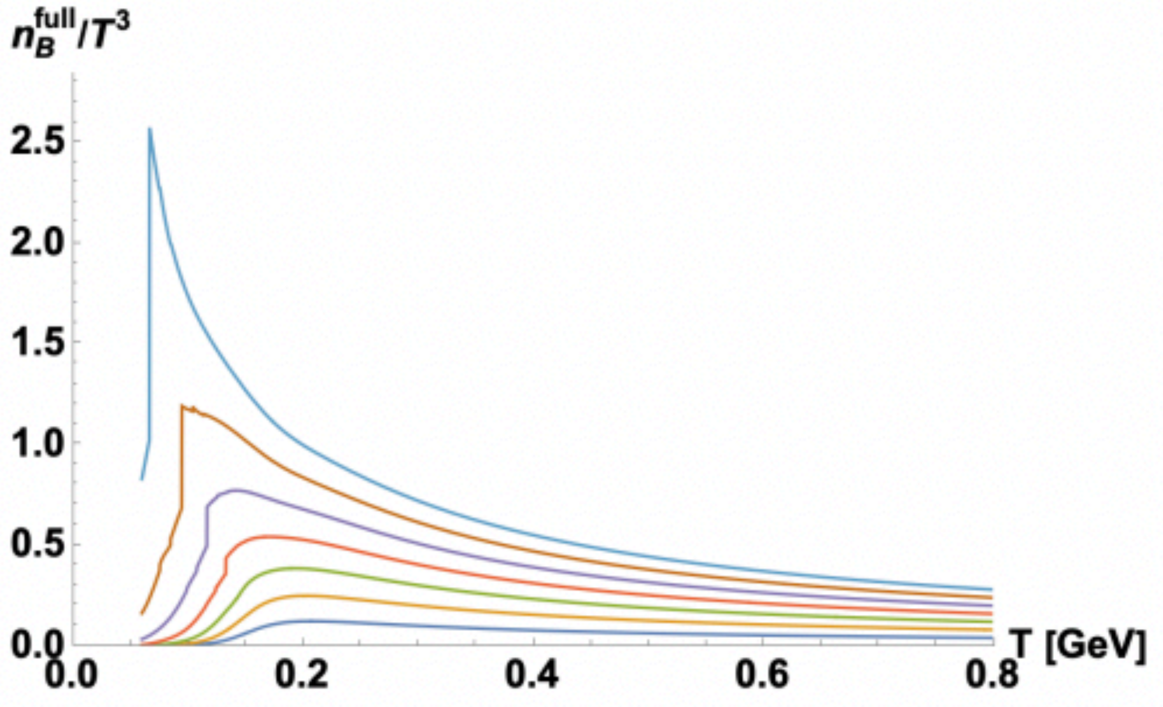
Introducing a Critical Point

$$T'_{crit}(T,\mu_B) \approx T_0 + \left(\left. \frac{\partial \chi^B_{2,lattice}(T,0)}{\partial T} \right|_{T=T_0} \right)^{-1} \frac{n_B^{crit}(T,\mu_B)}{T^3(\mu_B/T)} + \ldots$$

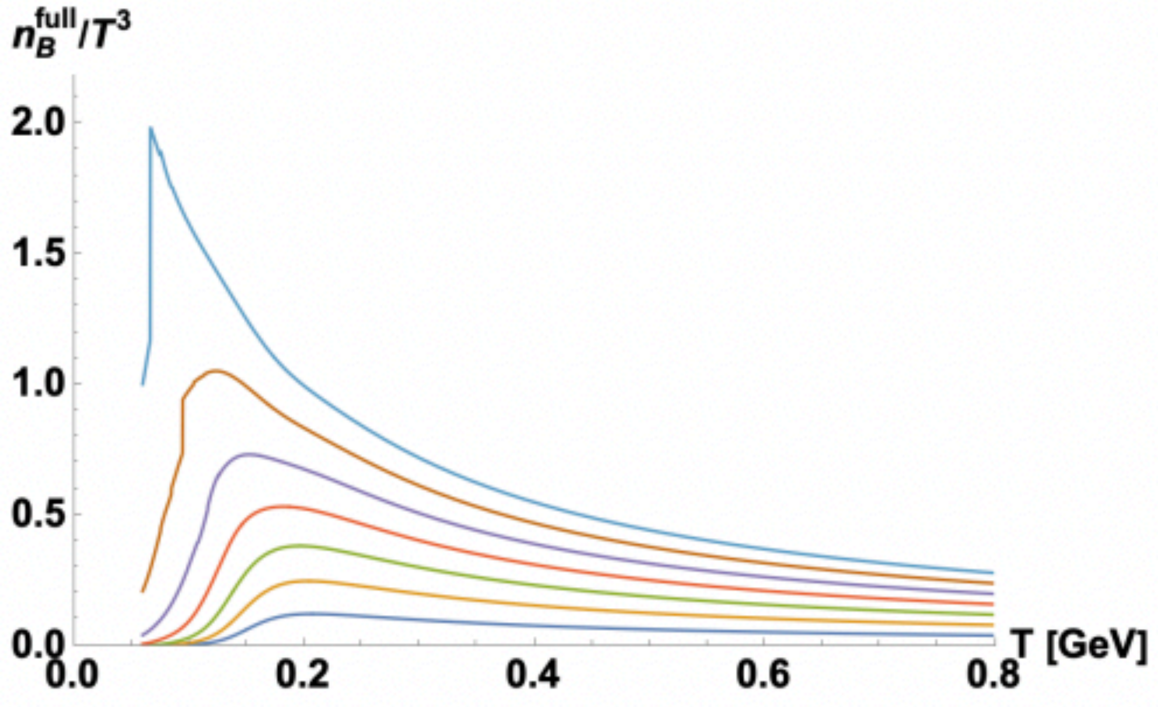
$$\chi_1^{crit}(T,\mu_B) = \frac{n_B^{crit}(T,\mu_B)}{T^3} = \frac{\partial(P^{crit}(T,\mu_B)/T^4)}{\partial(\mu_B/T)}$$

Baryon density results

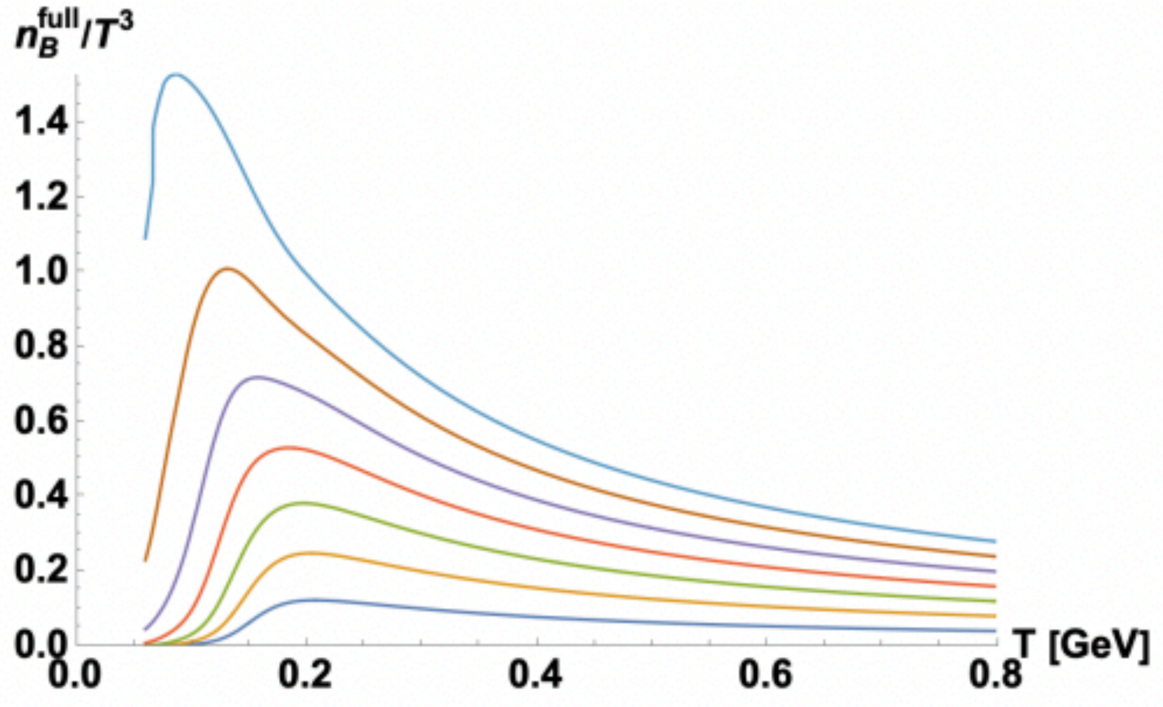
Full Baryon density at constant μ_B



$\mu_{BC} = 0.35GeV$

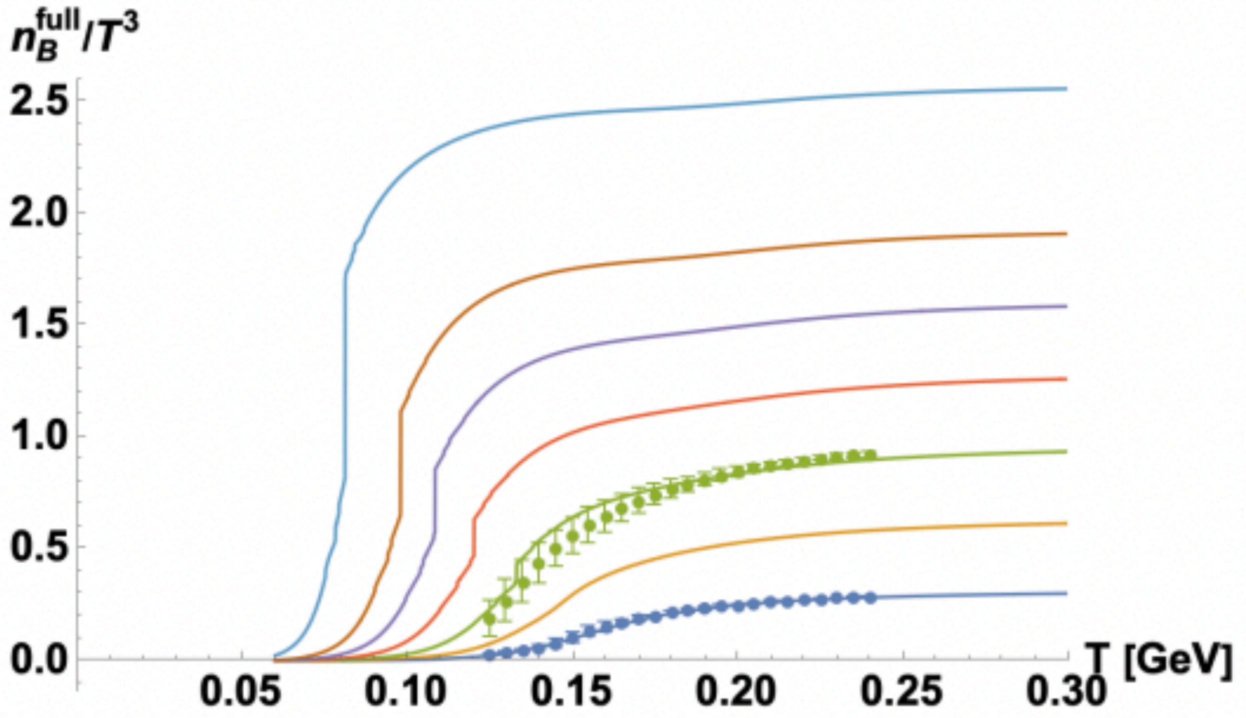


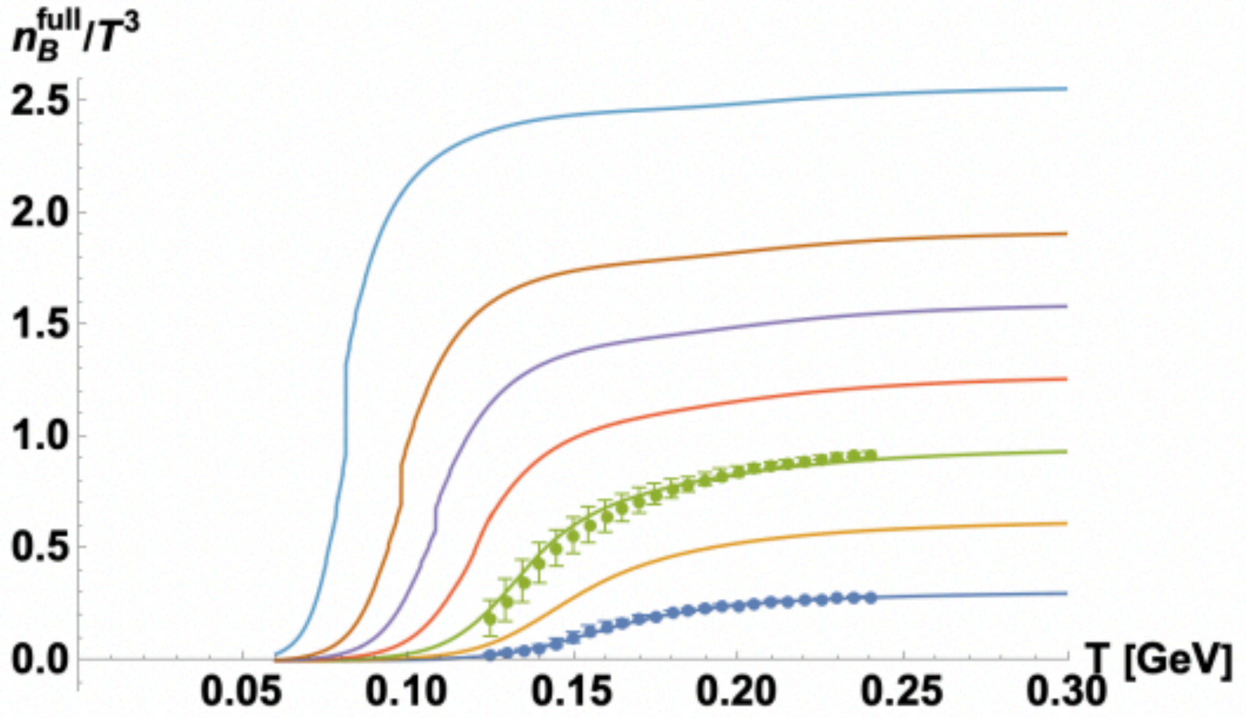
$\mu_{BC} = 0.55GeV$

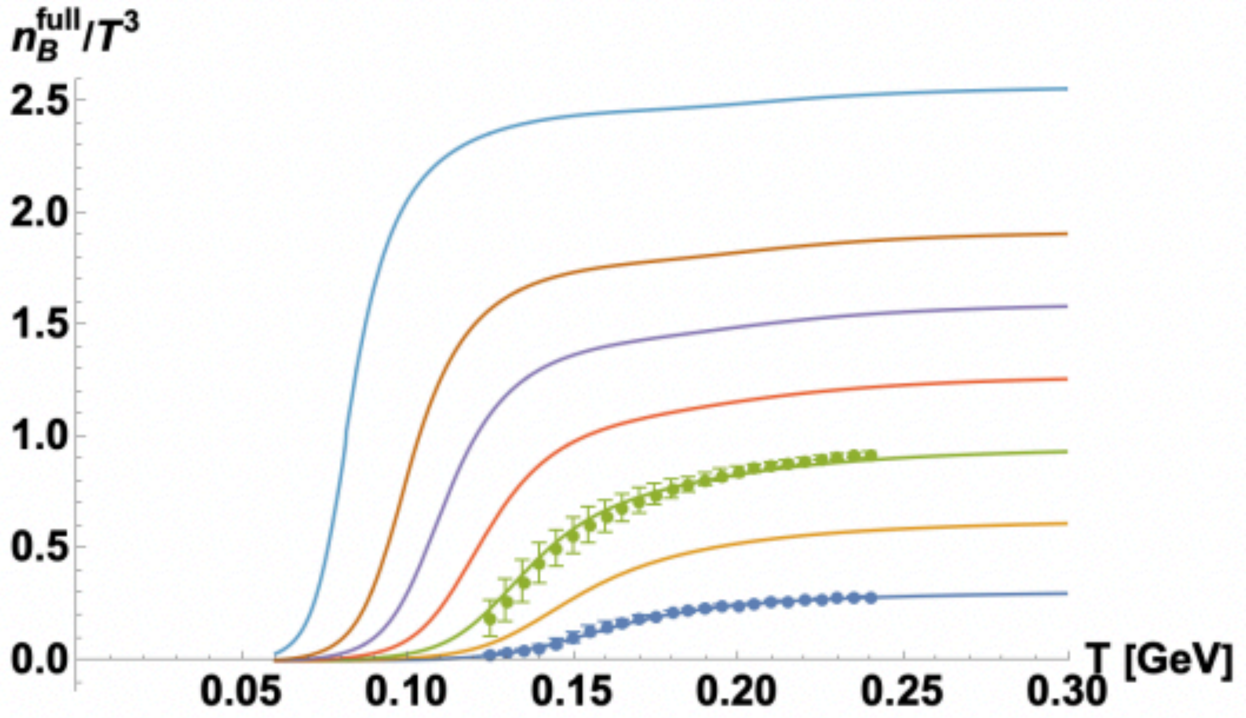


$\mu_{BC} = 0.65GeV$

Full Baryon density at constant $\frac{\mu_B}{T}$



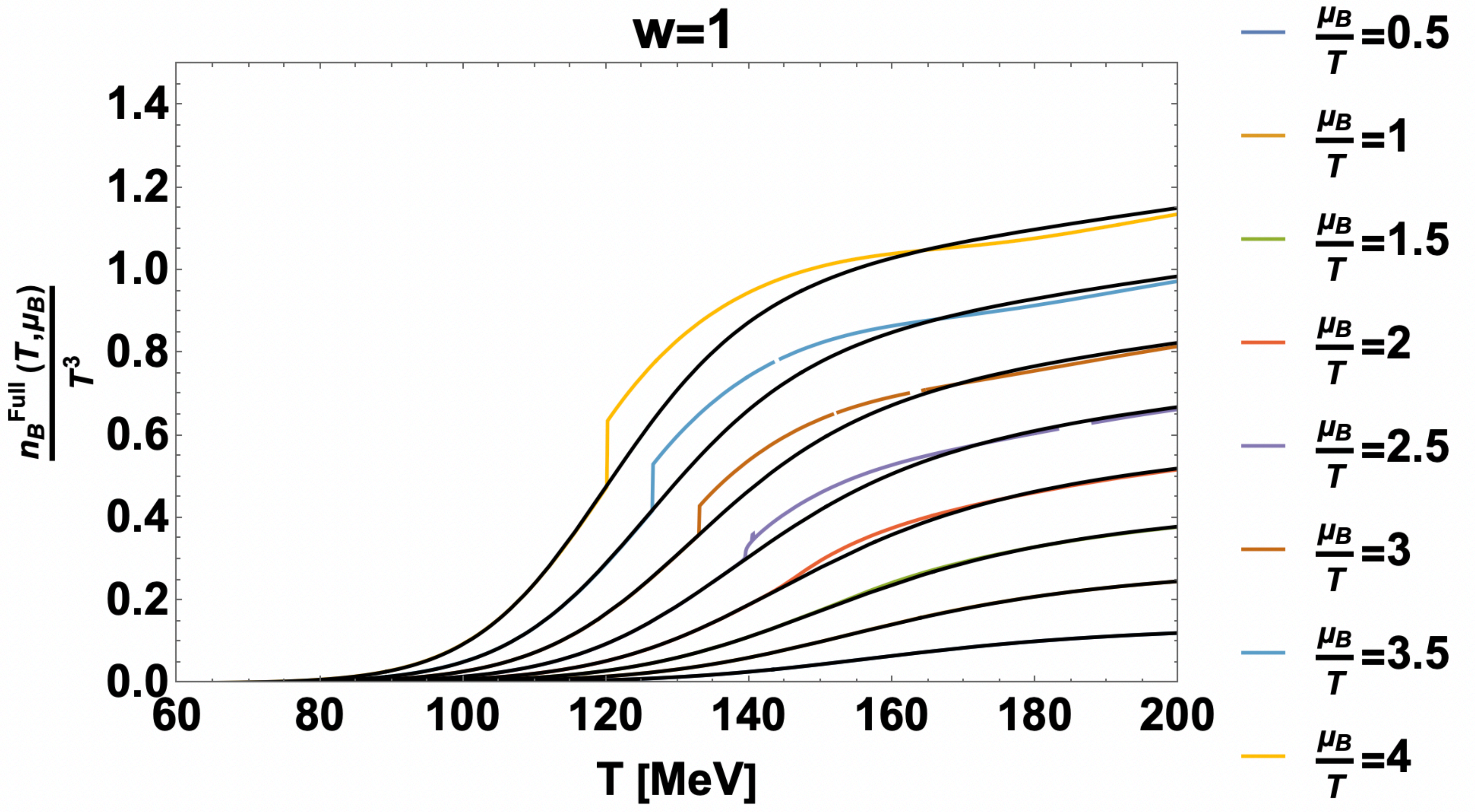


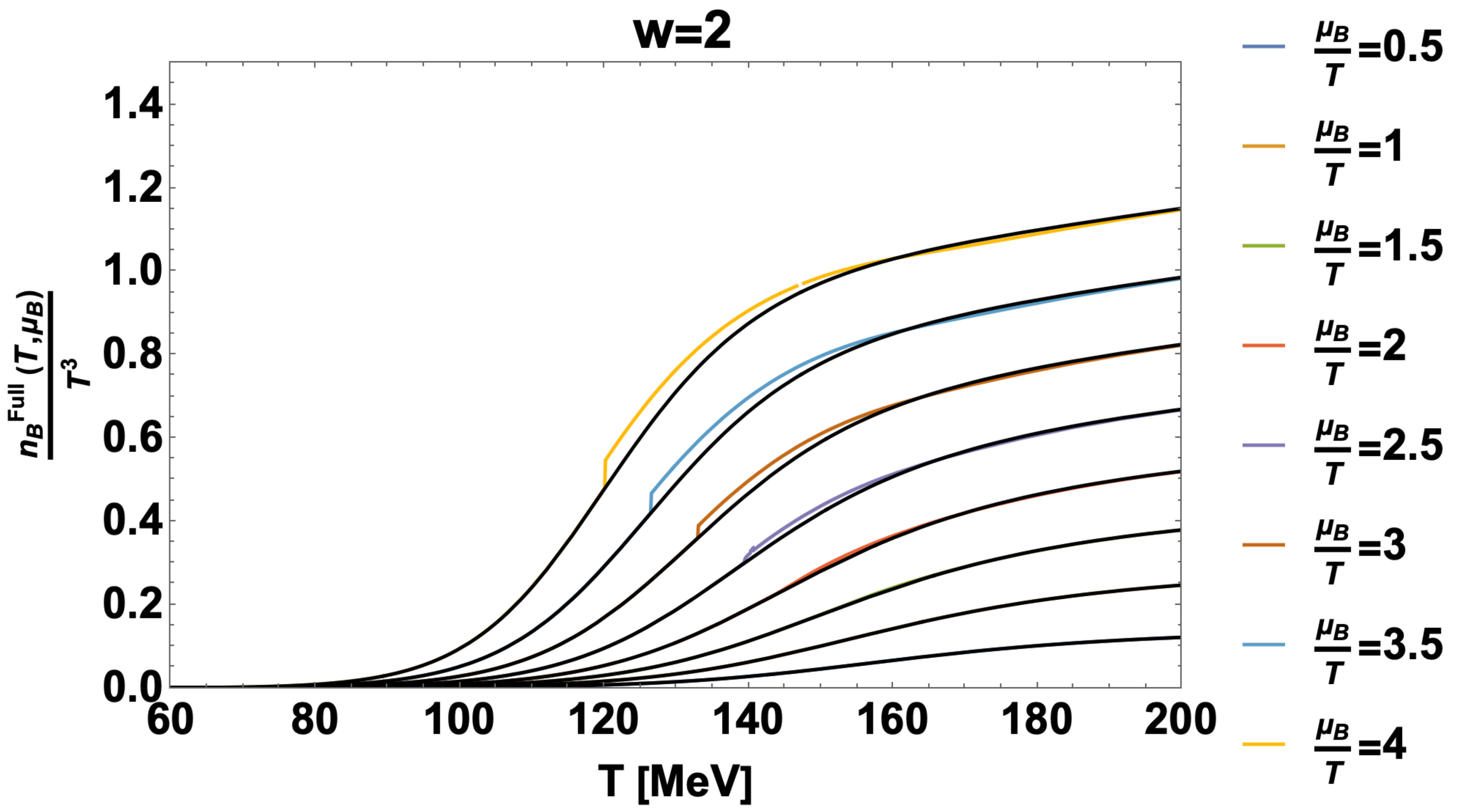


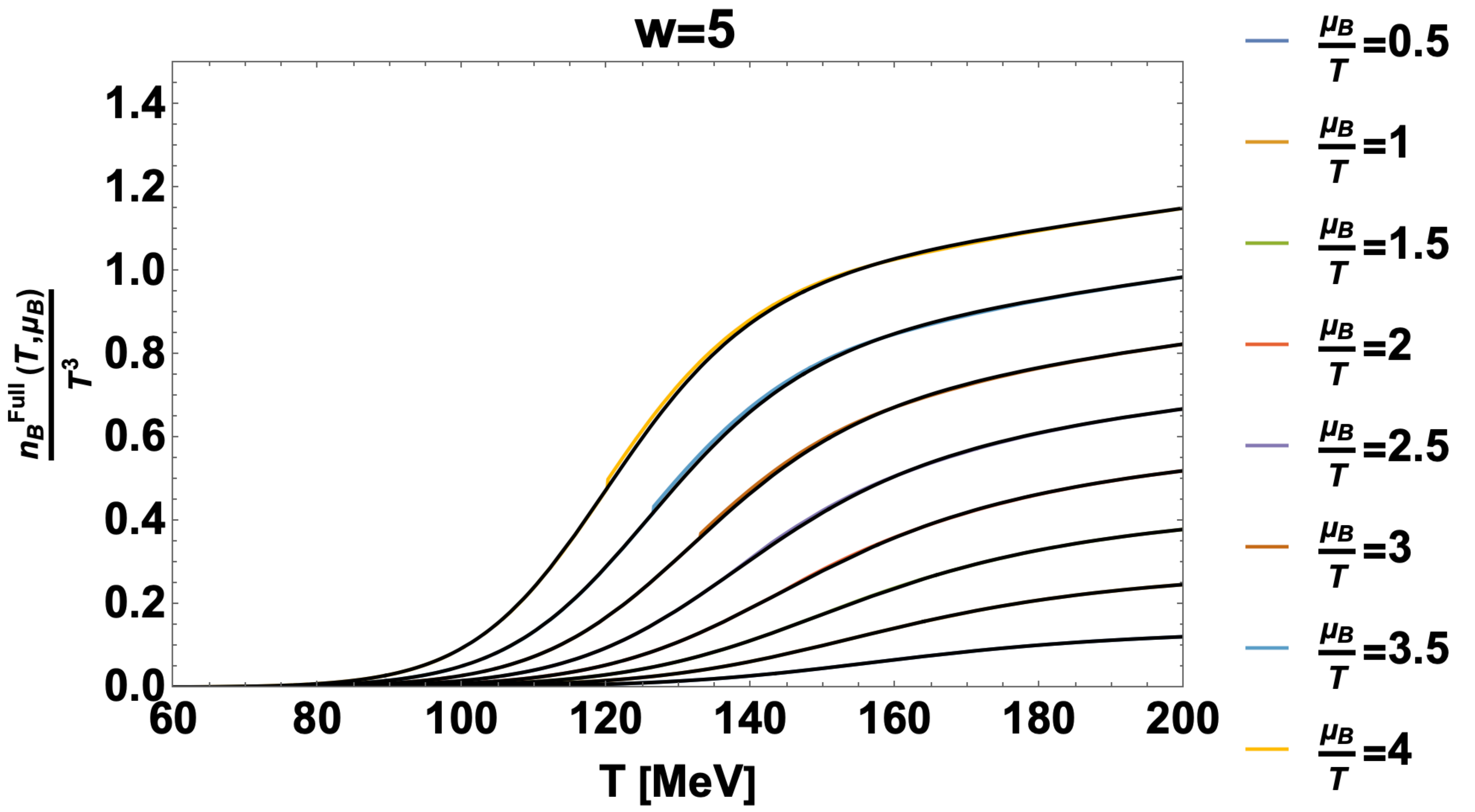
Estimating the Critical contribution

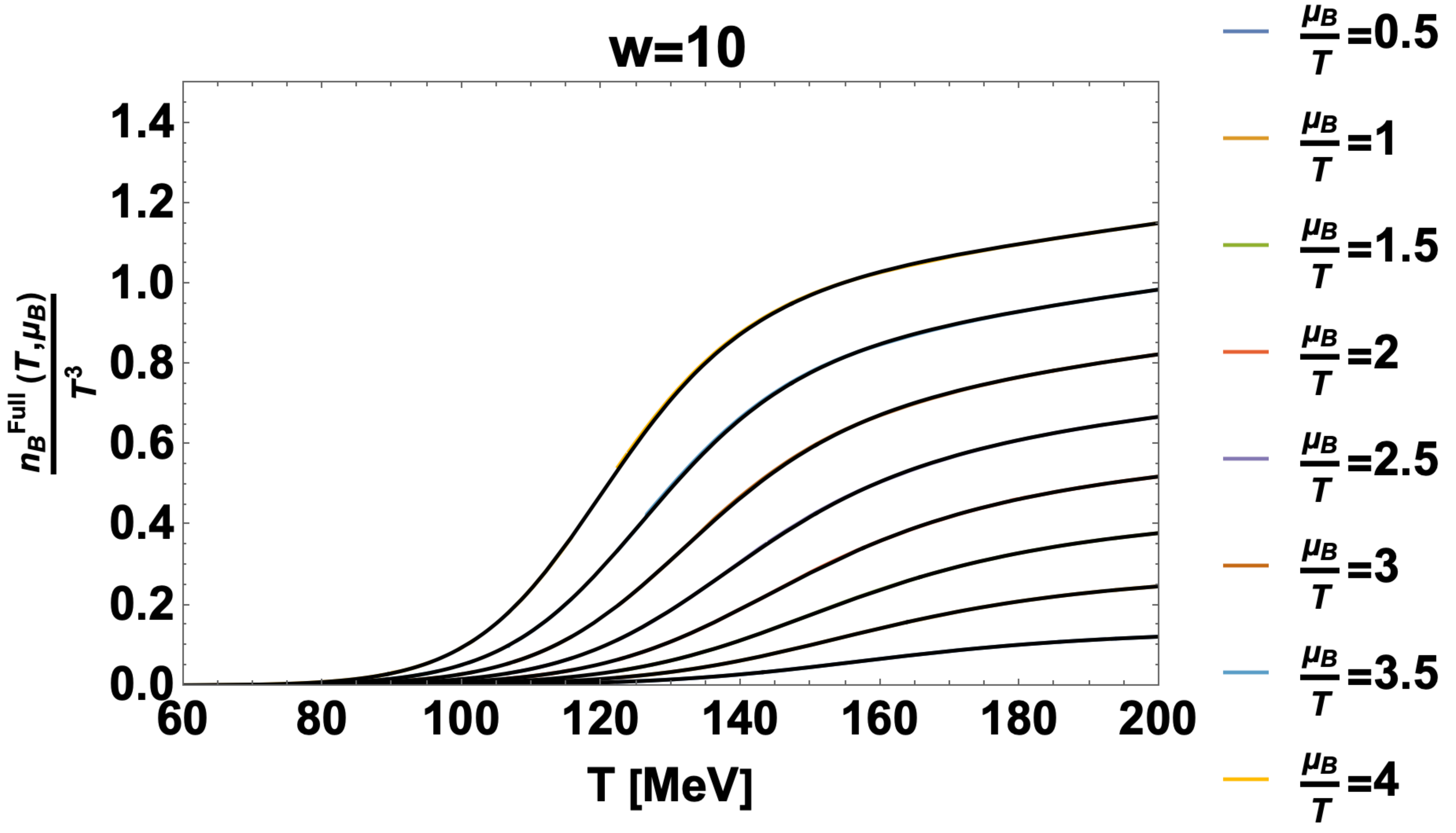
Baryon Density at a constant $\frac{\mu_B}{T}$ for w=1 to w=10

$\mu_B = 350\ [MeV],\ \ \alpha_{12} = 90,\ \ \rho = 2$





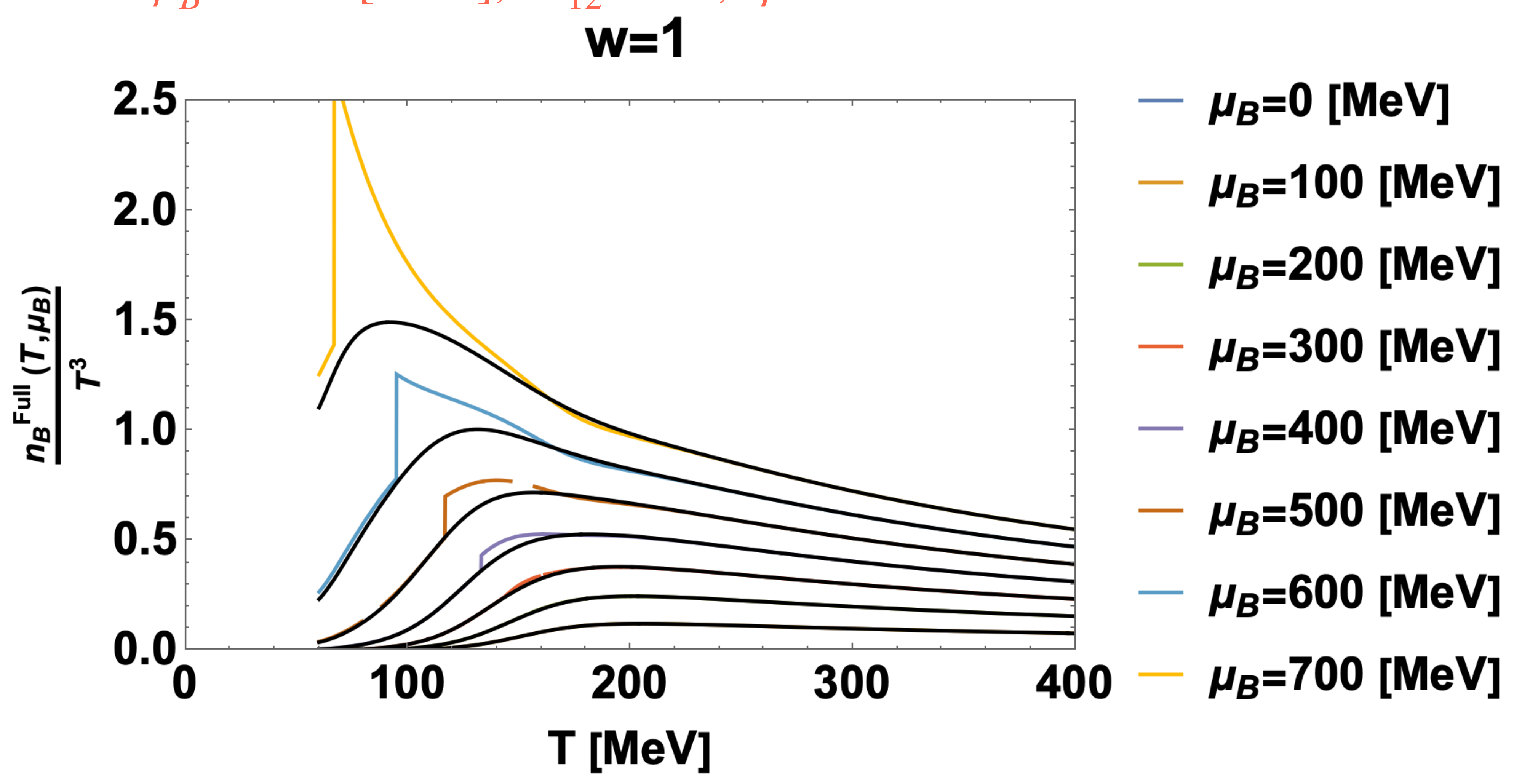


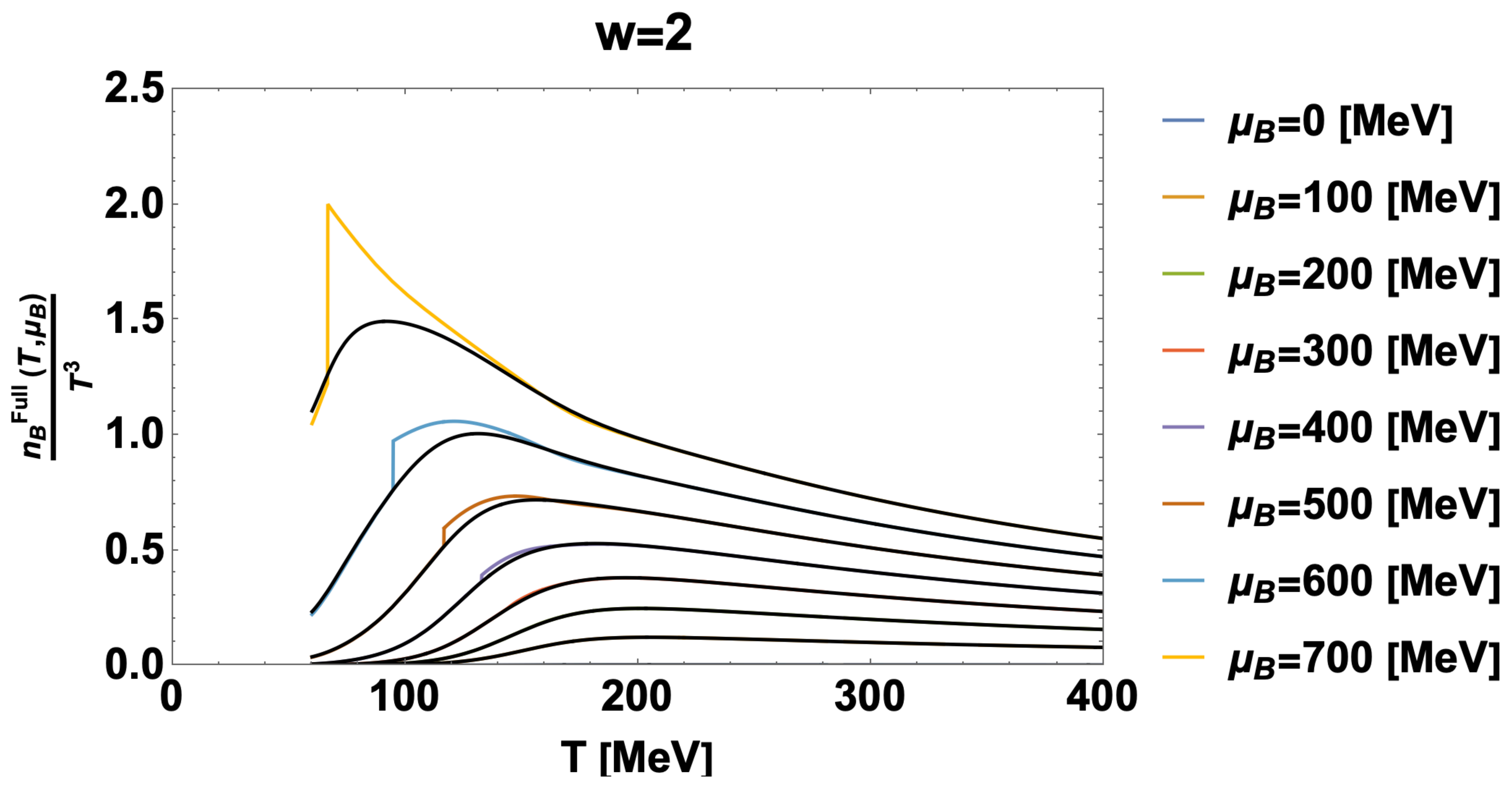


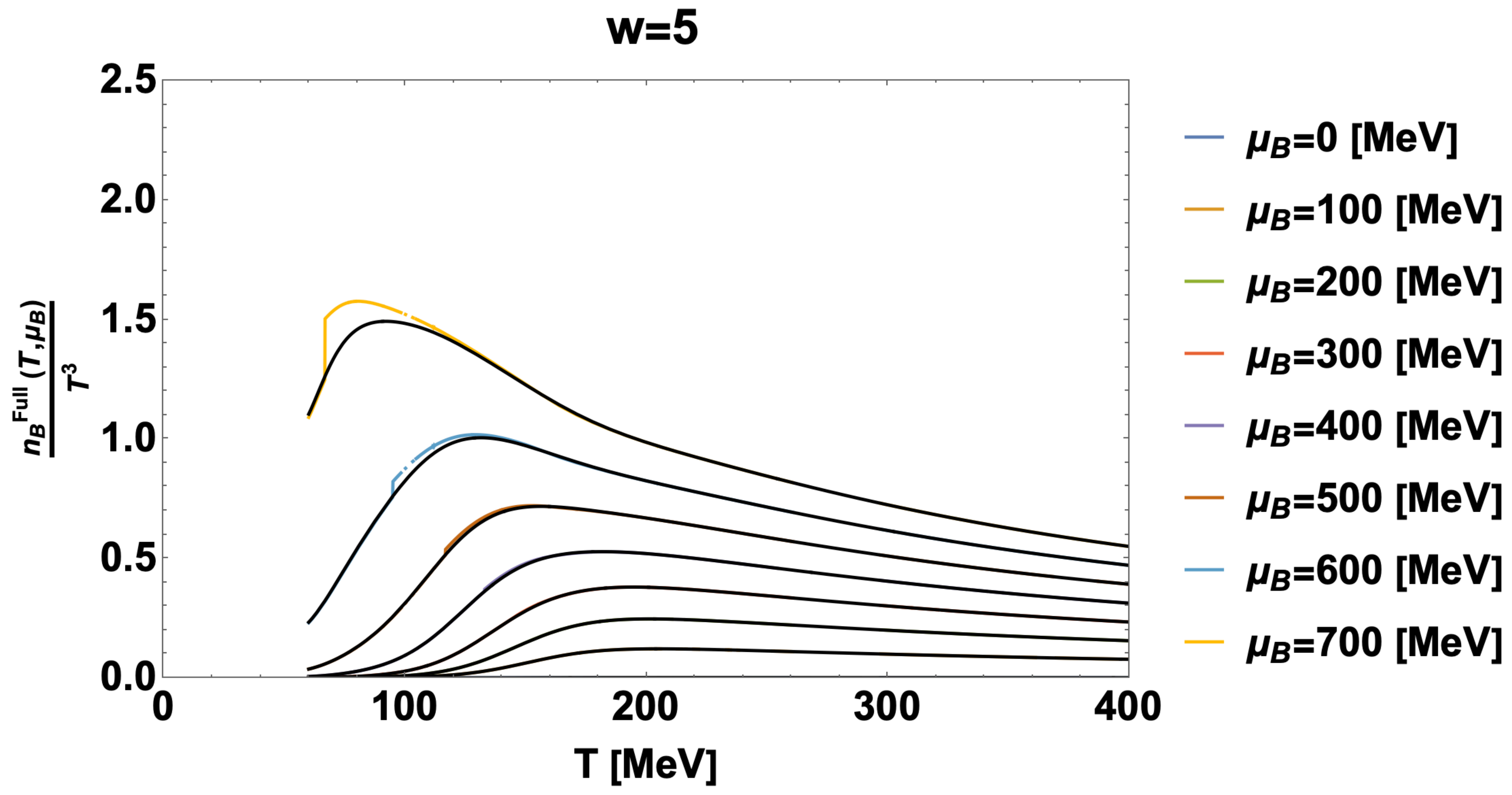
Estimating the Critical contribution

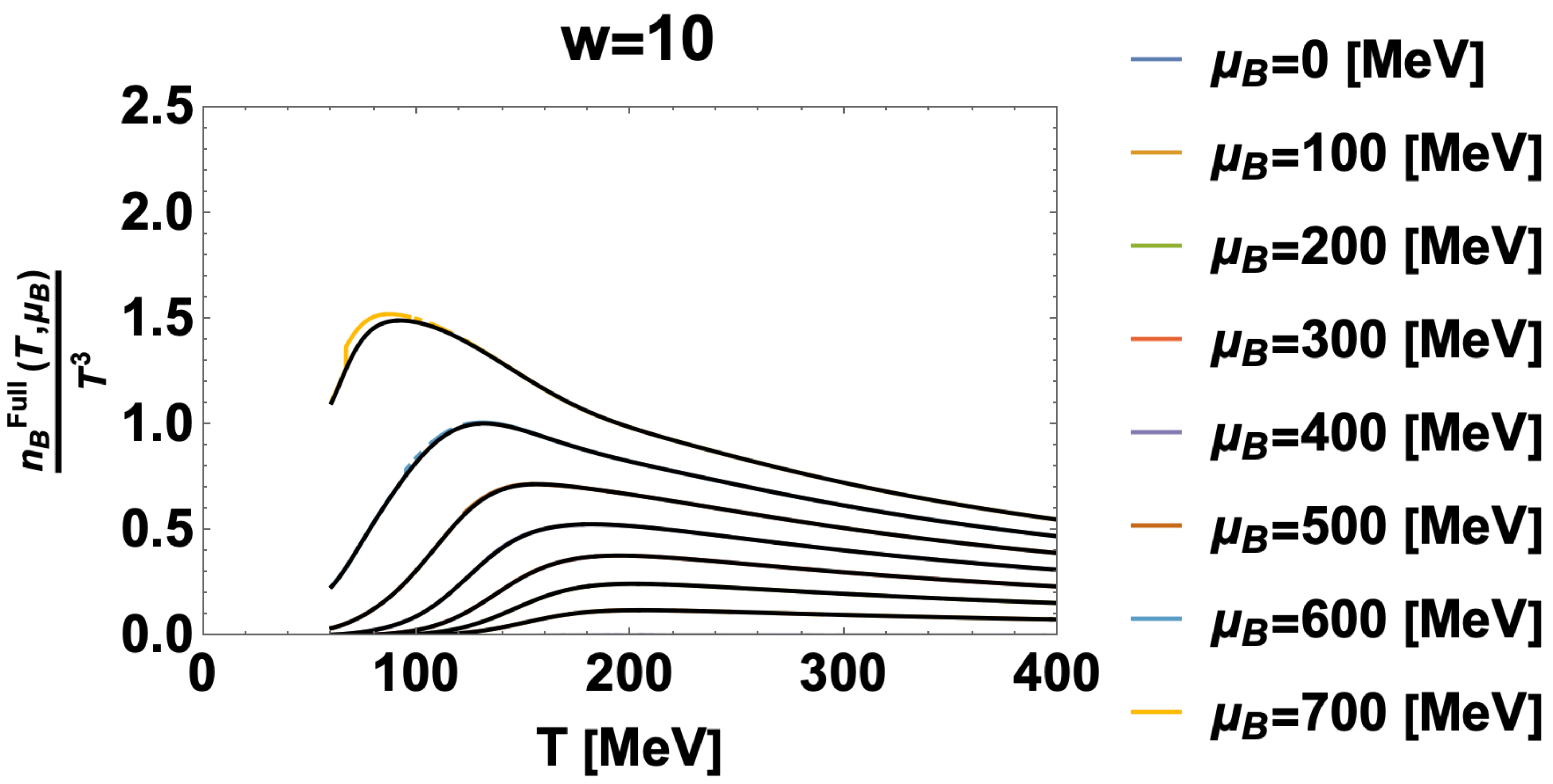
Baryon Density at a constant μ_B for 1 to w=10

$\mu_B = 350\ [MeV],\ \ \alpha_{12} = 90,\ \ \rho = 2$



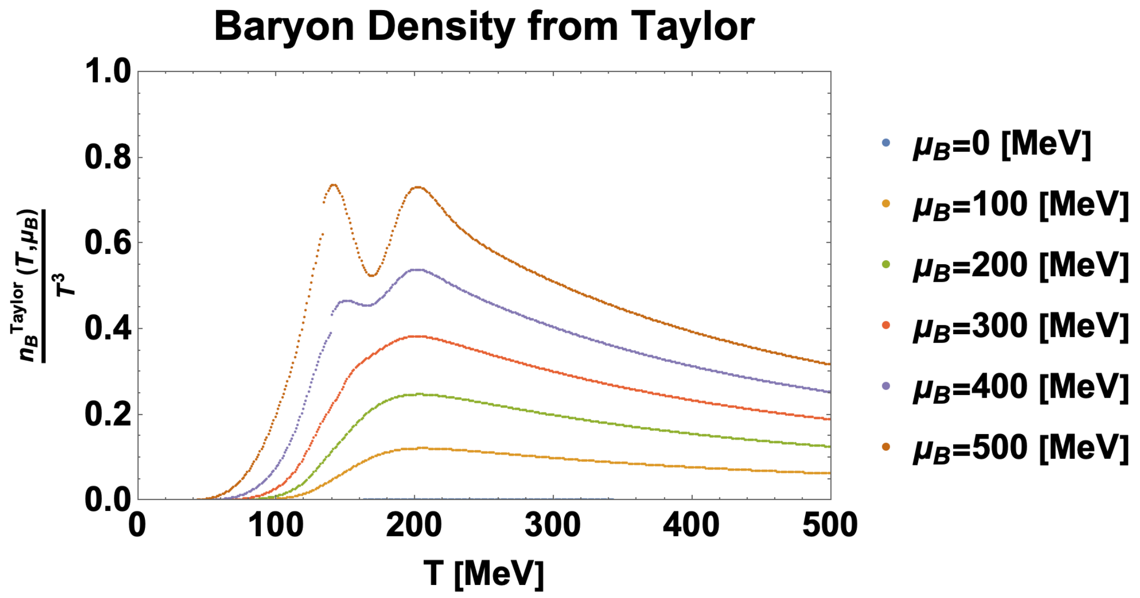


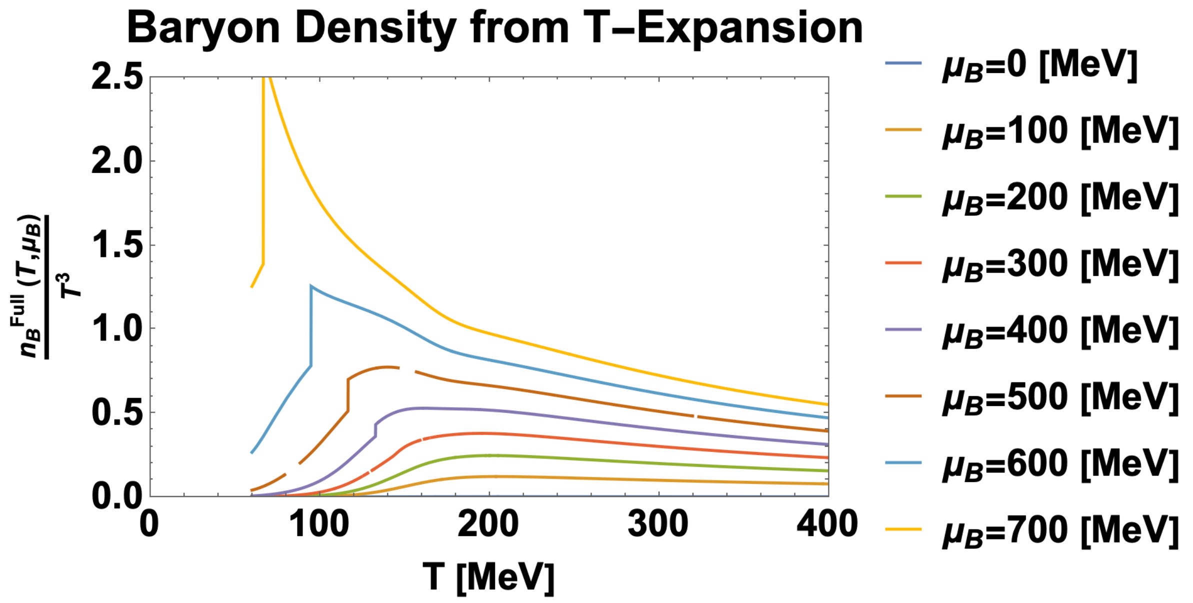


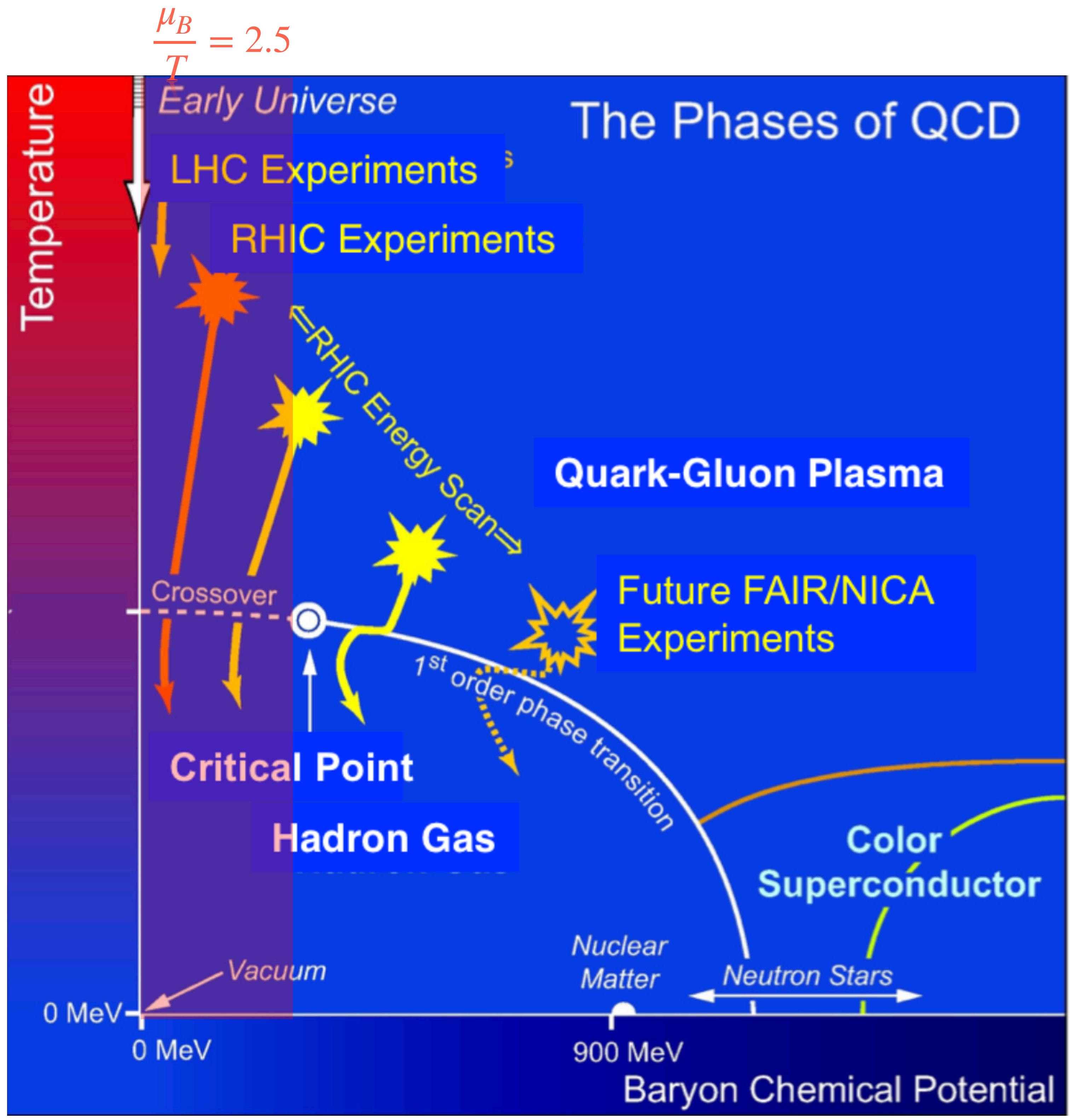


Summary

$\mu_B = 350\ [MeV],\ \ \alpha_{12} = 90,\ \ w = 1,\ \ \rho = 2$



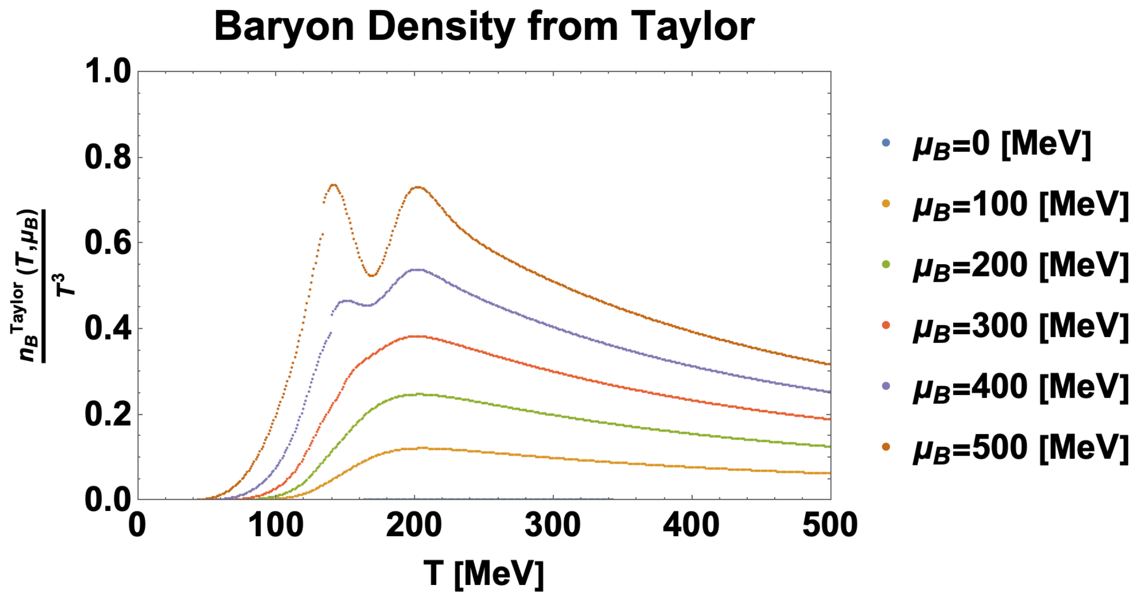


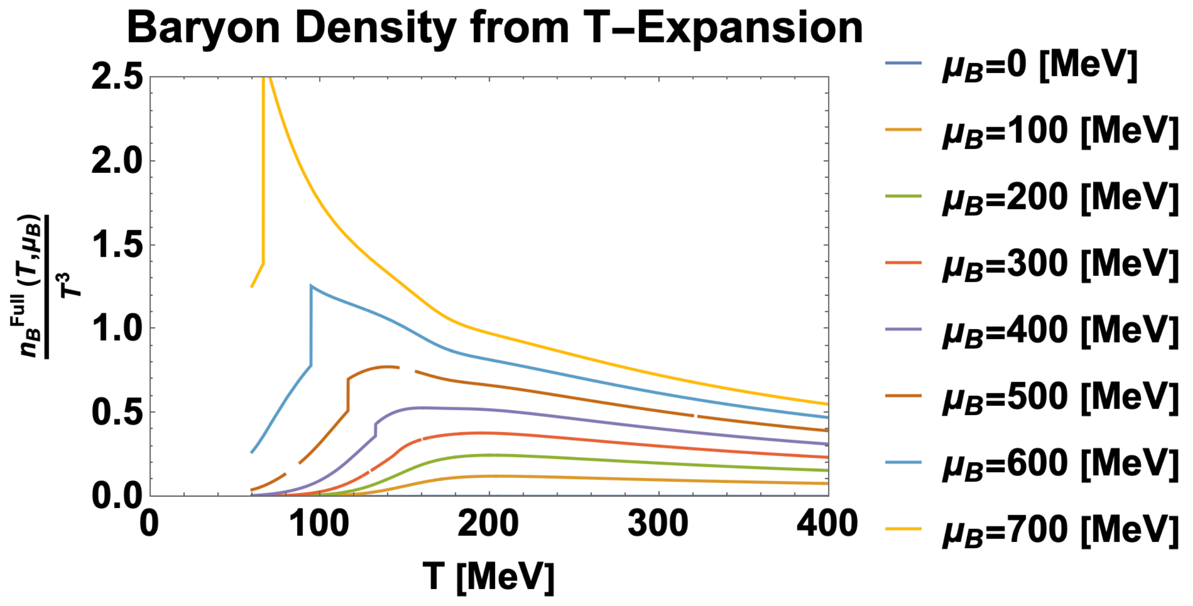


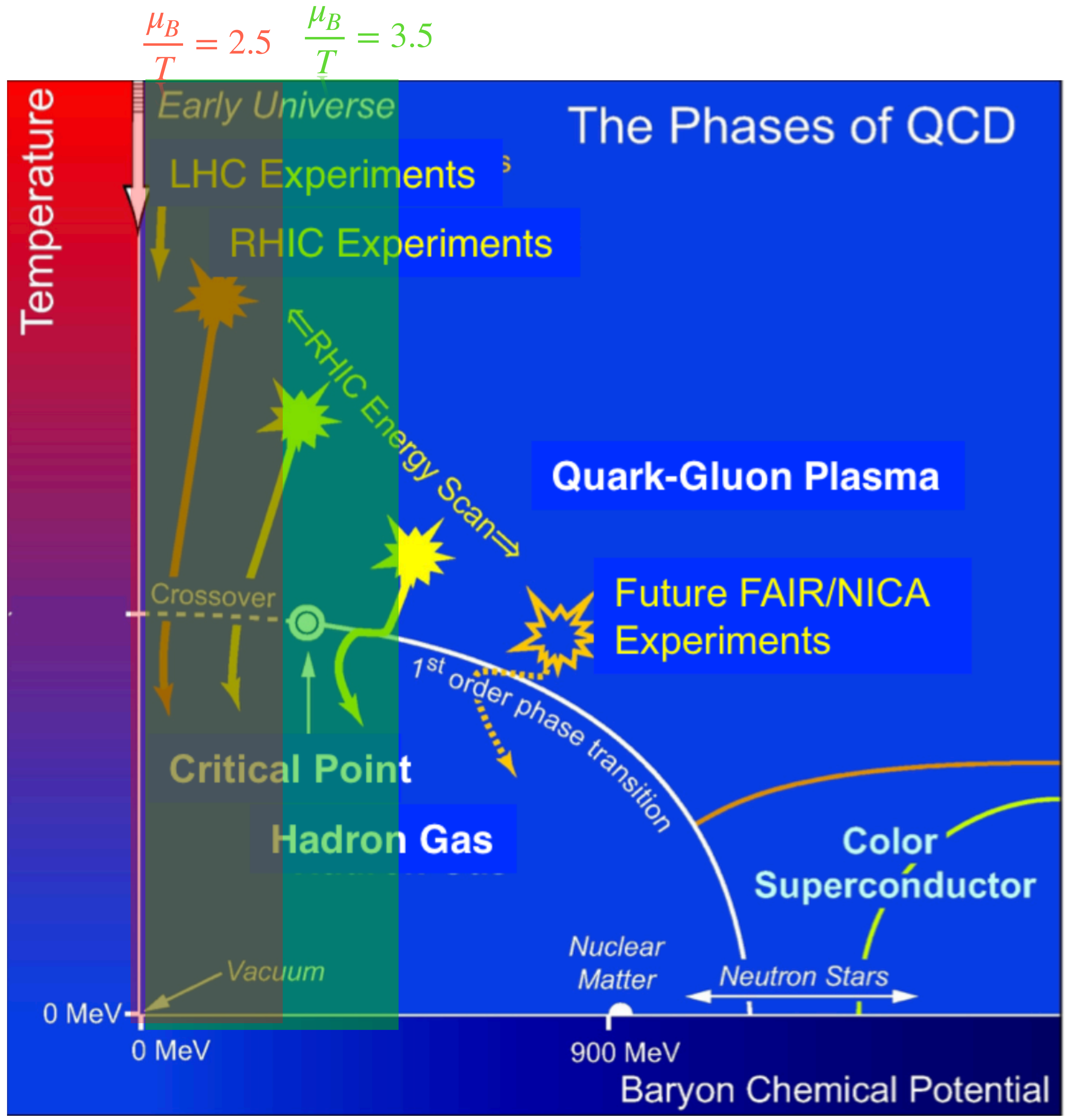
- **A more physical EoS that captures a large part of the Phase diagram is required.**
- **We provide a family of EoS with a correct Critical point up** $\mu_B = 700MeV$.
- **Our EoS allows users to change parameters and compare with the data from the Experiment (Beam Energy Scan II)**

Summary

$\mu_B = 350\ [MeV],\ \ \alpha_{12} = 90,\ \ w = 1,\ \ \rho = 2$







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Thank you for your attention!

Future Work

- **Compute all the Thermodynamic Observables (Pressure, Entropy, Energy density, Speed of Sound, etc)**

$$\frac{P(T,\mu_B)}{T^4} = \frac{P(T,0)}{T^4} + \int_0^\infty d\hat{\mu}'_B \frac{n_B(T,\hat{\mu}'_B)}{T^3}$$

$$\frac{S(T,\mu_B)}{T^3} = \frac{1}{T^3}\left(\frac{\partial P}{\partial T}\right)\bigg|_{\mu_B}$$

$$\frac{\epsilon(T,\mu_B)}{T^3} = \frac{S}{T^3} - \frac{P}{T^4} + \frac{\mu_B}{T}\frac{n_B}{T^3}$$

$$c_s^2(T,\mu_B) = \left(\frac{\partial P}{\partial \epsilon}\right)\bigg|_{S/n_B}$$

- **Explore and constrain the Parameters space by requesting thermodynamics Stability and causality of our EoS**

- **Merge with Nuclear Matter EoS at low Temperatures**